Stress analysis, buckling analysis and optimum proportions of an isotropic web-stiffened sandwich cylindrical shell under hydrostatic pressure.

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STRESS ANALYSIS, BUCKLING ANALYSIS AND OPTIMUM PROPORTIONS OF AN ISOTROPIC WEB-STIFFENED SANDWICH CYLINDRICAL SHELL UNDER HYDROSTATIC PRESSURE

Frank Robert Kroner
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by

FRANK ROBERT KRONER, JR.

B.S., United States Naval Academy (1961)

SUBMITTED IN PARTIAL FULFILLMENT OF THE REQUIREMENTS FOR THE DEGREES OF NAVAL ENGINEER AND MASTER OF SCIENCE IN OCEAN ENGINEERING

at the

MASSACHUSETTS INSTITUTE OF TECHNOLOGY

May, 1970
STRESS ANALYSIS, BUCKLING ANALYSIS AND OPTIMUM PROPORTIONS
OF AN ISOTROPIC WEB-STIFFENED SANDWICH CYLINDRICAL SHELL
UNDER HYDROSTATIC PRESSURE

by
Frank Robert Kroner, Jr.

Submitted to the Department of Naval Architecture and Marine
Engineering on May 21, 1970, in partial fulfillment of the
requirements for the degree of Naval Engineer in Naval
Architecture and Marine Engineering and Master of Science in
Ocean Engineering.

ABSTRACT

Web-stiffened sandwich cylindrical shells consisting of
two concentric cylindrical shells joined and stiffened by
annular rings (webs) were considered in the investigation.
The purpose of this investigation was to study the stress
distribution and buckling modes, and to determine the scant-
lings and configuration of optimum proportions.

The scantlings and configuration of optimum proportions
was defined as that with the lowest weight to displacement
ratio at a given loading of hydrostatic pressure.

The scantlings (outer shell thickness, inner shell
thickness, web thickness, web depth (separation of shells)
and configuration (web spacing) were systematically and
independently varied and the effect upon failure mode pres-
sures, and the weight to displacement ratio was determined.
The results were plotted graphically.

The form of stress and buckling analysis results follow
closely those of a ring-stiffened cylinder. Geometric sim-
ilitude was determined to exist. The scantlings and configu-
ration of optimum proportions were found to be a function of
the loadings.

Web-core sandwich configurations possess structural
efficiencies on the order of 20% higher than those of con-
ventional ring-stiffened construction for loadings on the
order of 1000 psi (2225 feet) and these diminish with depth
to about 5% higher for loadings on the order of 4000 psi
(8900 feet).

The real advantage of the web-stiffened sandwich con-
struction lies in the use of thinner plating with all its
superior metallurgical and fabrication qualities.

Thesis Supervisor: Professor J. H. Evans
Title: Professor of Naval Architecture
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NOTATION

b  Web thickness

\( (E_s/E_T) - 1 \)

C  Web constant

D  Bending, rigidity, \( Et^3/(1-v^2) \)

\( D_M \)  Mean diameter of shells

d  RRO-RRI

E  Young's modulus

\( E_s \)  Secant modulus

\( E_T \)  Tangent modulus

f  Stress ratio, \( \sigma_x/\sigma_\phi \)

\( g_o, g_i \)  Edge coefficients for outside and inside shells

\( g_{oo}, g_{oi}, g_{io}, g_{ii} \)  Lamé deflection coefficients for the annular webs

h  \( R_o-R_i \)

\( k_i \)  Stress intensity per pound of pressure \( (K_x = \sigma_x/p) \)

L  Unsupported length of cylinder

\( L_B \)  Length between bulkheads

m, n, k  Integers

\( N_x, N_\phi \)  Longitudinal and circumferential forces per unit length

x, \( \phi, r \)  Longitudinal, circumferential, and radial coordinates
<table>
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<td>$T_o, T_O$</td>
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<td>$T_i, T_I$</td>
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</table>
$u, v, w$ Longitudinal circumferential and radial deflections

$W_{os}$, $W_{is}$ Radial deflection of outside and inside shell

$W_{ow}$, $W_{iw}$ Radial deflection of the outer and inner surfaces of the web

$z$ Radial bending rigidity, $E/(1-\nu^2)$

$\varepsilon_x'\varepsilon_\phi'\varepsilon_r$ Longitudinal, circumferential, and radial strains

$\sigma_x'\sigma_\phi'\sigma_r$ Longitudinal, circumferential, and radial stresses

$\theta_o, \theta_i$ Shell flexibility parameters for outside and inside shells,
\[ \left[ 3(1-\nu^2) \right]^{1/4} \left[ L/(Rt) \right]^{1/2} \]

$\psi$ Bending stress coefficient,
\[ [(1-\nu^2)/3]^{1/2} \]

$\nu$ Poisson's ratio for isotropic materials

$\sigma_i$ Stress intensity, Huber-Hencky, Von Mises, $\left[ \sigma_1^2 + \sigma_2^2 - \sigma_1 \sigma_2 \right]^{1/2}$

$\sigma_Y$ Yield stress

$\sigma_{ijkl}$ ($\sigma_{iKL}$) Stress location

\[
\begin{align*}
i & = x, \phi, r & \text{direction} \\
j & = F, M & \text{shell (web-shell juncture, midbay)} \\
K & = O, I, W & \text{outside shell, inside shell, web} \\
L & = O, I, M & \text{fiber (outer, inner, membrane)} \\
(jiKL) & \text{ i.e., MXOI = midbay, longitudinal, outside shell, inner fiber}
\end{align*}
\]

$P$ Yield failure at shell midbay
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INTRODUCTION

One of the more promising structural concepts for the design and fabrication of cylindrical pressure hull structures is the sandwich concept. Structural engineers in the aircraft industry have long recognized and taken advantage of the favorable strength-weight characteristics of sandwich-type construction. In studying the literature, however, it has been found that the loading conditions encountered in these applications have dictated sandwich structural arrangements that would be of no direct use in the design of pressure hulls for submersibles.

The demands of hydrostatic pressure loading are such that in order to exploit the sandwich concept for pressure hull construction a core comprised of elements which are compression resistant as well as shear resistant is sought. The major difference is that membrane loads are predominant in hydrospace applications, whereas in aerospace, the bending and shear type loads are of prime concern.

Experimental programs and analytical studies have been going on concurrently in order to develop rational formulas based on thin-shell theory for predicting the static structural response of these types of structures. In reference 10 Pulos presents an analysis of the axisymmetric elastic
deformations and stresses in a web-stiffened sandwich cylinder under hydrostatic pressure. Raetz presents a similar analysis for the toroidal tube-stiffened sandwich cylinder.\textsuperscript{11} In references 12 and 14 Nott presents a simplified stress and strain analysis.

Nott's equations were used and programmed for the stress analysis of the web-stiffened structure.

In reference 22 the axisymmetric and asymmetric buckling equations are developed. In references 15 and 23 an expression for the general instability is determined.

The above equations are utilized, with others derived for web stability, to study the parametric nature of web-core sandwich structure and to determine optimum proportions.

Optimum design must be based on rational considerations of inelastic behavior, however, and not on the "one horse shay" concept based on elastic considerations of instability. The "ignorance factors" can then be representative of the variability introduced by certain intangibles which are not easily considered in a theory.

Some of the intangibles which influence static strength and complicate the problem so that appropriate design formulas cannot be derived on purely theoretical grounds are those inherent in the fabrication process itself:
(1) Initial stress.
(2) Residual stress.
(3) Imperfect circularity.
(4) Non-isotropic or non-homogeneity of material.
(5) Actual boundary conditions at stiffening rings.

Therefore, in view of these complications, the development of satisfactory design criteria must first be predicted on rigorous mathematical theory with its concomitant idealizations, and then, empirical factors derived from test data can be introduced to "adjust" the theories to take account of these many variables which could not or were not considered.

Thus an inelastic analysis of all instability modes is used.
GENERAL THEORY

1. Stress Analysis

A theoretical analysis of the axisymmetric elastic deformations and stresses in a web-stiffened sandwich cylindrical shell structure under external hydrostatic pressure is presented in Appendix A. The solution is based on the use of edge coefficients for plate and shell elements of finite length, and includes the computation of the edge forces and moments arising at the common junctions of these elements.

Equations are given for computing numerically the longitudinal and circumferential stresses in the two coaxial cylindrical shells and the radial and tangential stresses in the web stiffeners between the two shells.

To evaluate the structural strength of cylindrical sandwich shells, the locations and magnitudes of maximum stresses must be determined. Rigorous analyses for stress distributions in sandwich shells loaded under external hydrostatic pressure were carried out by Pulos and Raetz. These analyses illustrate that the maximum stresses occur in the inner and outer shells at locations next to the annular webs and midway between the two webs that separate the shells. Therefore, the stresses in the inside and outside shells at these two locations are of
interest for the purpose of structural design and evaluation.

The procedure and nomenclature for programming the stress analysis is that of Nott.12,14

a. NOMENCLATURE
b. ASSUMPTIONS

From symmetry considerations it is seen that the edges of the web stiffener, circular annulus, do not undergo any rotation. This stems from the fact that a horizontal tangent or zero-slope condition is assumed to exist at the junctions of the webs with the two cylindrical shells. This assumption implies that the edge moments on each shell at the shell-web juncture balance each other, so that there are no net moments to be resisted by the web. Further, it is assumed that the web elements do not take any axial force due to the axial pressure, but that this is all resisted by the cylindrical shells.

Also inherent in the analysis (see Appendix A) is the neglect of the beam-column effect due to the axial portion of the hydrostatic pressure.
\[
D \frac{d^4w}{dx^4} + N_x \frac{d^2w}{dx^2} + \frac{Et}{R^2} W = P - \frac{\nu}{R} N_x
\]

neglecting the beam-column effect becomes

\[
D \frac{d^4w}{dx^4} + \frac{Et}{R^2} W = P - \frac{\nu}{R} N_x
\]

Thus the analysis of the web stiffeners reduces to that of a circular annulus subjected to axisymmetric in-plane radial forces on both its inner and outer boundaries.

On the basis of these assumptions, it is only necessary to derive edge coefficients for an annulus undergoing radial deflection.

The functions \( M, N, P, \) and \( Q \) are the Lamé deflection coefficients for the annular webs and can be expressed

\[
M = \frac{R_w}{b} \left( 1 + \frac{d}{2R_w} \right) \left( \frac{2R_w}{d} + \frac{d}{2R_w} - 2\nu \right)
\]

\[
N = \frac{R_w}{b} \left( 1 + \frac{d}{2R_w} \right) \left( \frac{2R_w}{d} + \frac{d}{2R_w} - 2 \right)
\]

\[
P = \frac{R_w}{b} \left( 1 - \frac{d}{2R_w} \right) \left( \frac{2R_w}{d} + \frac{d}{2R_w} + 2 \right)
\]

\[
Q = \frac{R_w}{b} \left( 1 - \frac{d}{2R_w} \right) \left( \frac{2R_w}{d} + \frac{d}{2R_w} + 2\nu \right)
\]

and the functions \( g_o, g_i \) are the edge coefficients from the outside and inside shells:
\[
g_o = - \frac{2R_o^2}{T_o L} F_{10}
\]
\[
g_i = - \frac{2R_i^2}{T_i L} F_{11}
\]

where

\[
F_{10} = \frac{\theta_o}{2} \left[ \frac{\sinh \theta_o + \sin \theta_o}{\cosh \theta_o - \cos \theta_o} \right]
\]
\[
F_{11} = \frac{\theta_i}{2} \left[ \frac{\sinh \theta_i + \sin \theta_i}{\cosh \theta_i - \cos \theta_i} \right]
\]

\[
\theta_o = \sqrt{3(1-v^2)} \cdot \frac{L}{\sqrt{R_o T_o}}
\]
\[
\theta_i = \sqrt{3(1-v^2)} \cdot \frac{L}{\sqrt{R_i T_i}}
\]

Thus, from Appendix A, the axial and transverse forces \((V_o, V_i, H_o, \text{ and } H_i)\), which act at the intersections of the annular webs and cylinders can be determined.

\[
V_i = \frac{T_i}{T_o} V_o
\]
\[
V_o = \frac{R_o/2}{1 + \frac{T_i}{T_o} \cdot \frac{R_i}{R_o}}
\]
\[(g_0 - M) H_0 - NH_i = -\frac{R_0^2}{T_o} \left( 1 - \frac{\nu/2}{1 + \frac{T_i}{T_o} \cdot \frac{R_i}{R_0}} \right) + \frac{b}{2} M \]

\[-P H_0 + (g_i - Q) H_i = \frac{\nu}{2} \frac{R_0^2}{T_i} \left( 1 - \frac{1}{1 + \frac{T_i}{T_o} \cdot \frac{R_i}{R_0}} \right) + \frac{b}{2} P \]

And from these forces the maximum stresses in the shells can be determined.

\[P \]

\[\text{c. MAXIMUM STRESSES} \]

\[\text{OUTER SHELL} \]

1. \[\sigma_{xfoo} = \left[ + \frac{V_o}{R_o} - \frac{H_o}{L} \right. \left. (B_{40}) \right] \sigma_{uo} \]

\[\sigma_{\phi foo} = \left[ + 1.0 + \frac{H_o}{L}(B_{50}) \right] \sigma_{uo} \]
2. \( \sigma_{xfoi} = \left[ + \frac{V_0}{R_0} + \frac{H_0}{L} (B_{40}) \right] \sigma_{uo} \)

\( \sigma_{\phi foi} = \left[ + 1.0 - \frac{H_0}{L} (B_{60}) \right] \sigma_{uo} \)

3. \( \sigma_{xmoi} = \left[ + \frac{V_0}{R_0} + \frac{H_0}{L} (B_{30}) \right] \sigma_{uo} \)

\( \sigma_{\phi moi} = \left[ + 1.0 - \frac{H_0}{L} (B_{30}) \right] \sigma_{uo} \)

4. \( \sigma_{xmoi} = \left[ + \frac{V_0}{R_0} - \frac{H_0}{L} (B_{30}) \right] \sigma_{uo} \)

\( \sigma_{\phi moi} = \left[ + 1.0 - \frac{H_0}{L} (B_{20}) \right] \sigma_{uo} \)

where \( \sigma_{uo} = - p \frac{R_0}{T_0} \)

**INNER SHELL**

5. \( \sigma_{xfoi} = \left[ + \frac{V_i}{R_i} - \frac{H_i}{L} (B_{4i}) \right] \sigma_{ui} \)

\( \sigma_{\phi fio} = \left[ - \frac{H_i}{L} (B_{5i}) \right] \sigma_{ui} \)

6. \( \sigma_{xfi} = \left[ + \frac{V_i}{R_i} + \frac{H_i}{L} (B_{4i}) \right] \sigma_{ui} \)

\( \sigma_{\phi fii} = \left[ - \frac{H_i}{L} (B_{6i}) \right] \sigma_{ui} \)
7. \( \sigma_{xmi0} = \left( + \frac{V_i}{R_i} + \frac{H_i}{L} (B_{3i}) \right) \sigma_{ui} \)

\( \sigma_{phi0} = \left( - \frac{H_i}{L} (B_{3i}) \right) \sigma_{ui} \)

8. \( \sigma_{xmi1} = \left( + \frac{V_i}{R_i} - \frac{H_i}{L} (B_{3i}) \right) \sigma_{ui} \)

\( \sigma_{phi1} = \left( - \frac{H_i}{L} (B_{2i}) \right) \sigma_{ui} \)

where \( \sigma_{ui} = -p \frac{R_i}{T_i} \)

and

\( B_{10} = 2F_{20} - \nu F_{30} \)

\( B_{20} = 2F_{20} + \nu F_{30} \)

\( B_{30} = F_{30} \)

\( B_{40} = F_{40} \)

\( B_{50} = 2F_{10} + \nu F_{40} \)

\( B_{60} = 2F_{10} - \nu F_{40} \)

\( F_{10} = \frac{\theta_o}{2} \left( \frac{\sinh \theta_o + \sin \theta_o}{\cosh \theta_o - \cos \theta_o} \right) \)

\( F_{20} = \theta_o \left( \frac{\cosh \theta_o/2 \sin \theta_o/2 + \sinh \theta_o/2 \cos \theta_o/2}{\cosh \theta_o - \cos \theta_o} \right) \)

\( F_{30} = \frac{6\theta_o}{\sqrt{3(1-\nu^2)}} \left( \frac{\cosh \theta_o/2 \sin \theta_o/2 - \sinh \theta_o/2 \cos \theta_o/2}{\cosh \theta_o - \cos \theta_o} \right) \)
\[ F_{40} = \frac{3\theta_o}{\sqrt{3(1-v^2)}} \frac{\sinh \theta_o - \sin \theta_o}{\cosh \theta_o - \cos \theta_o} \]

\[ \theta_o = \frac{4}{\sqrt{3(1-v^2)}} \cdot \frac{L}{\sqrt{R_O T_O}} \]

and likewise for \( B_{1i}, B_{2i}, \ldots, F_{4i}, \theta_i \) by substituting the inside \( I \) for the outside \( O \).

d. WEB STRESSES

From Appendix A the maximum stress in the web

\[ \sigma_{r\omega} = -p \frac{\left( \frac{-HHI(\sqrt{R})^2 + HHO}{b/2 (1-(\sqrt{R})^2)} \right)}{\left( \frac{-HHI(\sqrt{R})^2 + HHO(\sqrt{R})}{b/2 (1-(\sqrt{R})^2)} \right)} \]

\[ \sigma_{\phi\omega} = -p \frac{\left( \frac{-HHI(\sqrt{R})^2 + HHO}{b/2 (1-(\sqrt{R})^2)} \right)}{\left( \frac{-HHI(\sqrt{R})^2 + HHO(\sqrt{R})}{b/2 (1-(\sqrt{R})^2)} \right)} \]

where

\[ HHO = H_O + b/2 \]
\[ HHI = H_i \]
\[ RR = R_{RI}/R_{RO} \]
\[ RRI = R_i - T_i/2 \]
\[ RRO = R_O + T_O/2 \]
2. Buckling Analysis and Failure Modes

The axisymmetric and asymmetric (lobar) buckling equations developed by Lunchick$^4$ and Reynolds$^8$ for ring-stiffened cylinders are modified in Appendix B for application to web-stiffened sandwich cylinders. They are evaluated for both the inner and outer shells. Reference 22 indicates good agreement between theoretical and experimental collapse pressures.

The general instability equation used has found considerable experimental verification.$^{15,23}$

The two modes of web buckling are first, a variation of the Foppl or Levy ring buckling (effectively TOKUGAWA) and second, the buckling of an annulus in biaxial edge compression (see Appendix B).

a. FAILURE MODES

1. STRENGTH

Using the Huber-Hencky-Von Mises criterion

$$
\sigma_i = \sqrt{\sigma_{\phi}^2 + \sigma_X^2 - \sigma_{\phi} \sigma_X}
$$

$$
\sigma_i = P \sqrt{K_{\phi}^2 + K_X^2 - K_{\phi} K_X}
$$

$$
\therefore \quad p = \frac{\sigma_y}{\sqrt{K_{\phi}^2 + K_X^2 - K_{\phi} K_X}}
$$

evaluated at the web-shell interface (PF), midbay (P), and outer fiber of web (PWS).
2. AXISYMMETRIC BUCKLING

\[ \text{PPC1} = \frac{E_T L^2}{12(1-\nu^2)K_x \pi^2 R^2} \left[ \frac{\pi^2 R^2 t^2 (1+\frac{3c}{4})}{L^4} + \frac{12(1+c)(1-\nu^2)}{1+\frac{3c}{4}} \right] \]

where

\[ c \equiv \frac{E_s}{E_T} - 1 \]
\[ \nu \equiv \frac{1}{2} - \frac{E_s}{E(1/2-\nu_E)} \]

3. ASYMMETRIC BUCKLING

\[ \text{PPC2} = \frac{\pi^2 E_T f_E t}{3K_x \phi (1-\nu^2)} \frac{t}{R} \left[ \frac{(R_t)^{1/2}}{L} \right]^2 \frac{1+\frac{3\phi c}{4}}{3-2\phi (1-f_E)} \]

where

\[ \phi = 1.23 \left( \frac{R_t}{L} \right)^{1/2} \]
\[ f_E \equiv \frac{K_x}{K_\phi} \]

4. GENERAL INSTABILITY

\[ \text{PPC3} = \sqrt{E_T \cdot E_s} \cdot \frac{t_T}{R} \left[ \frac{\lambda^4}{n^2 - l + \frac{\lambda^2}{2}} \cdot \frac{(n^2 + \lambda^2)(1-3/4)}{l} \cdot \frac{1-E_s}{E_T} \cdot \frac{\lambda^4}{n^2 + \lambda^2 (n^2 + \lambda^2)^2} \right] \]
\[ + \frac{(n^2 - l) E_T I_{EFF}}{R^3 L} \]
where

\[ \lambda = \frac{\pi R}{L_B} \]

\[ I_{\text{eff}} \]

for

\[ t_T = T_0 + t_i \]

\[ R = R_w \]

5. **WEB BUCKLING**

a. **Ring Buckling**

\[ \text{PPC4} = \frac{24 E_T I_{\text{eff}}}{D_m^3 L (1-\nu^2)} \]

b. **Edge Compression**

\[ \text{PPC5} = \frac{4\pi^2 E_T}{12 (1-\nu^2) (K_\varphi+4K_R)} \left( \frac{b}{h} \right)^2 \]
\[ I_{\text{eff}}: \]

\[ I_C = \frac{ba^3}{12} + \frac{LT_o^3}{12} + \frac{LT_i^3}{12} + T_o L \left( \frac{d}{2} - \frac{T_o}{2} \right)^2 \]

\[ + T_i L \left( \frac{d}{2} - \frac{T_i}{2} \right)^2 \]

\[ A = (T_o L + T_i L + db) \]

\[ y = (Y_{cg} - \frac{d}{2}) \]

\[ Y_{cg} = \frac{T_o L \left( d - \frac{T_o}{2} \right) + db \left( \frac{d}{2} \right) + T_i L \left( \frac{T_i}{2} \right)}{(T_o L + db + T_i L)} \]

\[ I_{\text{eff}} = I_C - A y^2 \]
AXISYMMETRIC BUCKLING (BETWEEN WEBS)

ASYMMETRIC (LOBAR) BUCKLING (BETWEEN WEBS)

GENERAL INSTABILITY (CONCURRENT FAILURE OF SHELLS AND WEBS)

FAILURE MODES (SHELL)
WEB (RING) BUCKLING

WEB (COMPRESSION) BUCKLING

FAILURE MODES (WEB)
PROCEDURE

The buckling pressure is a function of the cylinder geometry and the secant and tangent moduli as determined from a stress-strain intensity diagram for the shell material.

Before the buckling equations can be used, $E_s$ and $E_T$ must be related to the applied pressure. The secant and tangent moduli are defined

$$E_s = \frac{\sigma_i}{\varepsilon_i}$$

$$E_T = \frac{d\sigma_i}{d\varepsilon_i}$$

and are shown graphically:
Hence $E_s$ and $E_T$ are readily determined from the stress-strain curve (see Appendix C). Having determined $E_s/E$ and $E_T/E$ as functions of $\sigma_1$, they must be applied to the hydrostatically loaded cylindrical shell. By expressing hydrostatic pressure in terms of the stress intensity, a relationship between $E_s$, $E_T$, and pressure will be established.

\begin{align*}
(1) \quad P_c &= f(\text{geometry}, E_s, E_T) \\
\text{a buckling equation} \\
(2) \quad \sigma_i &= \sqrt{\sigma_1^2 + \sigma_2^2 - \sigma_1 \sigma_2} = p\sqrt{K_1^2 + K_2^2 - K_1 K_2} \\
\text{a stress intensity (Huber-Hencky-Von Mises) equation}
\end{align*}

The intersection of these two curves will predict failure by inelastic buckling if below the yield strength.

Thus the general solution technique is to (1) solve for the state of stress, and consequently a relationship
between $\sigma_i$ and $p$, and then (b) solve the buckling modes as a function of $\sigma_i(E_s,E_T)$. The intersection of the curves being the critical pressure (see lb).

1. Solution Technique.

a. STRESS ANALYSIS

INPUT DATA:

$R_o, R_I, T_o, T_I, b, L, \nu, p$

Let $p = 1.0$ psi

(1) Formulate $\theta_o, \theta_I$

(2) Formulate $F_{10}, F_{11}, F_{20}, F_{21} \ldots$

(3) Formulate $g_o, g_i$

(4) Formulate $R_w, d$

(5) Formulate $M, N, P, Q$

(6) Formulate $V_o, V_I$

(7) Solve simultaneous equations for $H_o, H_i$

(8) Formulate $HHO, HHI, RR$

(9) Formulate $\sigma_{uo}, \sigma_{ui}$

(10) Formulate $B_{10}, B_{11}, B_{20}, B_{21} \ldots$

(11) Solve for the states of stress
b. FAILURE MODES

(1) INELASTIC MATERIALS

(2) ELASTIC - PERFECTLY PLASTIC MATERIAL
c. **OPTIMUM PROPORTIONS**

A simple optimization program was used to determine optimum proportions. As indicated in the flow diagram, a comparison is made between the design collapse pressure (PCO) and all the failure mode pressures ($P_i$). The failure mode pressures consist of:

- $P = \text{yield at midbay}$
- $PF = \text{yield at web-shell junction}$
- $PWS = \text{yield in web}$
- $PPC_1 = \text{axisymmetric buckling, each shell}$
- $PPC_2 = \text{asymmetric buckling, each shell}$
- $PPC_3 = \text{general instability}$
- $PPC_4 = \text{web instability (ring)}$
- $PPC_5 = \text{web buckling (edge compression)}$
OPTIMIZATION PROGRAM FOR WEB-STIFFENED SANDWICH PRESSURE HULL

DIMENSION RD(5),TD(5),TI(5),B(5),WL(5)
COMMON NUE,ET,EZ,SIGMAY,LH
REAL NUE,LT,EFF,KR,KP

61 READ(5,100) PCO
100 FORMAT(F10.0)
     IF(PCO-1.0)3,3,60
3 CALL EXIT
60 READ(5,101) RHO,ET,EZ,SIGMAY,NUE,LB,R1,(RD(I),I=1,5)
101 FORMAT(F10.2,E10.2,F10.2,F10.0,F10.2,5F10.2)
     WRITE(6,200) PCO,RHO,NUE,SIGMAY,R1
200 FORMAT(' C Collapsed Pressure=',F10.2,'/5X',', Densitiy=',F10.0,'/5X',', Width=',
     F10.2,'/5X',', SIGMAY=',F10.0,'/5X',', Inside radius=S=',F10.2)
     WRITE(6,250)
250 FORMAT(5X,' RRO=',4X,' TTU=',4X,' TTI=',4X,' HR=',4X,' HR=',4X,' HST=',4X,' RST=',4X,' PR=',4X,' PR=',4X)
     WB=10.00
     PI=3.1416
     DO 25 J=1,5
     CO 20 K=1,5
       CO 10 M=1,5
     CO 5 I=1,5
     RRO=RD(I)
     BB=B(L)
     TTU=TD(J)
       TTI=TI(K)
     WWL=WL(M)
     RRI=R1
     CALL STAS(RRO,RRI,TTU,TTI,BD,WWL,SX,SPI,FX,FPHI,SSX,
       SSPHI,FFX,FFPHI,HR,HR,HTS,HTS,HS,HTS,EFF,D,H,HTS)
     RW=(RHO+RRI)/2.0
     CM=2.0*RA
     CALL FMCS(SX,SPI,FX,FPHI,HR,HR,SSX,SSPHI,FFX,FFPHI,HR,HR,HTS,HTS,HS,HTS,EFF,D,H,HTS)

32
1     H, Ts, BB, P, PF, PPC1, PPC2, PPC3, PPC4, PPC5, PNS, PND
  IF (PGO-PF) 4, 4, 7
  IF (PGO-P) 6, 6, 7
  IF (PGO-PPC1) 8, 8, 7
  IF (PGO-PPC2) 9, 9, 7
  IF (PGO-PPC3) 11, 11, 7
  IF (PGO-PPC4) 12, 12, 7
  IF (PGO-PPC5) 33, 33, 7
  IF (PGO-PNS) 34, 34, 7
  7    GO TO 5
  34   CALL FMCS (SSX, SSPH1, FEX, FEXH1, KR, KP, RR1, ITI, IEFF, 04, 0)
       IF (PGO-PF) 2, 2, 14
       IF (PGO-P) 13, 13, 14
       IF (PGO-PPC1) 16, 16, 14
       IF (PGO-PPC2) 17, 17, 14
       IF (PGO-PPC3) 18, 18, 14
       IF (PGO-PPC4) 19, 19, 14
       IF (PGO-PPC5) 21, 21, 14
       IF (PGO-PNS) 22, 22, 14
       14   GO TO 5
  22   WD=2.0*XRHL/64.4*(TTJ/RRC+ (RR1/RRU)*(TTI/RRC) + (RR1/RRU)* (RR1/RRU)
       1   * (0/(W WL+R5)))
       WRITE (6, 400) RRC, TTJ, ITJ, HRN, HT, HS, HTS, PF, WD
       IF (WD=WH) 30, 30, 31
  30   WB=WD
  31   GO TO 5
  5    CONTINUE
  10   CONTINUE
  15   CONTINUE
  20   CONTINUE
  25   CONTINUE
       WRITE (6, 300) WH
  300   FORMAT (' WEIGHT/SUJANCY=' F10.3 )
SUML.ROUTINE STAS(RG,RI,TC,TL,B,ML,MCQ,MPHI1U,FX1I1)
1 FPHI01,FXII,MPHI01,FX10,MPHI01,KR,KP,MRK,HTS,AS,HTS
2 JEFF,DO,HTS
COMMON NUE,E,T,E,Z,SIGMA,Y,LB
STRESS ANALYSIS (STAS) OF WEB STIFFENED SANDWICH CYLINDRICAL SHELLS
REAL NUE,L,E,N,MCQ,MPHI01,MCQ1,MPHI11,ML1,MPHIIU,
1 MXII,MPHI11,IEFF,KR,KP
PCP=1.0
L=XL
THEDA0=(3.0*(1.0-NUE**2.0))**0.25*L/(R*RI)**0.50
THEDA1=(3.0*(1.0-NUE**2.0))**0.25*L/(RI**1.1)**0.50
F10=THEDA0/2.0*(SINH(THEDA0)+SIN(THEDA0))/(COS(THEDA0))-(COS(THEDA0))
F11=THEDA1/2.0*(SINH(THEDA1)+SIN(THEDA1))/(COS(THEDA1))-(SIN(THEDA1))

BE0A0=THEDA0/2.0
BE0A1=THEDA1/2.0
F20=THEDA0*(COSH(THEDA0)-SIN(THEDA0)+SINH(THEDA0)COS(THEDA0))/
1 1
F21=THEDA1*(COSH(THEDA1)-SIN(THEDA1)+SINH(THEDA1)COS(THEDA1))/
1 1
F30=0.0*THEDA0/(3.0*(1.0-NUE**2.0))**0.5*(COSH(THEDA0)-SIN(THEDA0)-
1 1
SINH(THEDA0)COS(THEDA0))/(COSH(THEDA0)-SIN(THEDA0))
F31=0.0*THEDA1/(3.0*(1.0-NUE**2.0))**0.5*(COSH(THEDA1)-SIN(THEDA1)-
1 1
SINH(THEDA1)COS(THEDA1))/(COSH(THEDA1)-SIN(THEDA1))
F40=3.0*THEDA0/(3.0*(1.0-NUE**2.0))**0.5*(SINH(THEDA0)-SIN(THEDA0)
1 1
)/(COSH(THEDA0)-SIN(THEDA0))
F41=3.0*THEDA1/(3.0*(1.0-NUE**2.0))**0.5*(SINH(THEDA1)-SIN(THEDA1)
1 1
)/(COSH(THEDA1)-SIN(THEDA1))

GJ=-2.0*R0**2.0*F10/(10*L)
G1=-2.0*R1**2.0*F11/(11*L)
R=RI+RI
D=(RI-R1)*(1.0+1.0)/2.0
R=RW/0*(1.0+0.0)(2.0**R1)**(2.0**R1)/B**(2.0**R1)-2.0**NUE
N=RW/0*(1.0+0.0)**(2.0**R1)**(2.0**R1)/B**(2.0**R1)-2.0)
P=RW/0*(1.0+0.0)**(2.0**R1)**(2.0**R1)/B**(2.0**R1)+2.0)
\[ Q = RW/B \times (1.0 - B/2.0) \times (2.0 - RW)/D + D/2.0 \times RW + \beta \times NUT \] 
\[ VC = (R/R2 \times C)/(1.0 + (TI/TU) \times (RI/RU)) \] 
\[ VI = (TI/TU) \times VC \] 
\[ H2 = (((Q - Q) \times (R/R2 \times 2.0) \times (1.0 - (NUE/2.0)) + (1.0 - R/RI) \times (RI/RU)) + \beta \times MI \] 
\[ H1 = ((-P) \times (-P/R/U \times 2.0) \times (1.0 - (NUE/2.0)) + (1.0 - R/RI) \times (RI/RU)) + \beta \times MI \] 
\[ SIGUJ = -PCP \times RU/TU \] 
\[ SIGUI = -PCP \times RI/TI \] 
\[ BLJ = 2.0 \times F20 - NUE \times F30 \] 
\[ BIJ = 2.0 \times F20 - NUE \times F31 \] 
\[ B20 = 2.0 \times F20 + NUE \times F30 \] 
\[ B21 = 2.0 \times F21 + NUE \times F31 \] 
\[ B30 = F30 \] 
\[ B31 = F31 \] 
\[ B40 = F40 \] 
\[ B50 = 2.0 \times F40 - NUE \times F40 \] 
\[ B60 = 2.0 \times F40 - NUE \times F41 \] 
\[ B41 = F41 \] 
\[ B51 = 2.0 \times F41 + NUE \times F41 \] 
\[ B61 = 2.0 \times F41 + NUE \times F41 \] 
\[ C \text{ STRESS STATE ONE} \] 
\[ MX00 = (V0/R0 + H4 \times B0/L0) \times SIGU \] 
\[ MP10 = (V0/B0 \times H0/L0) \times SIGU \] 
\[ EX0 = (V0/R0 + H4 \times B0/L0) \times SIGU \] 
\[ F0 = 1.0 \times (1.0 + H0/B0/L0) \times SIGU \] 
\[ C \text{ STRESS STATE TWO} \] 
\[ EX10 = (V1/R1 + H4 \times B1/L1) \times SIGU \] 
\[ EP11 = (H1 \times B1/L1) \times SIGU \] 
\[ MX11 = (V1/R1 + H4 \times B1/L1) \times SIGU \] 
\[ MP11 = (H1 \times B1/L1) \times SIGU \] 
\[ C \text{ STRESS CALCULATION} \] 
\[ F11 = H1 + B/2.0 \]
DR1=41-T1/2.0
DR0=Ru+T0/2.0
RR=DR1/DR0
KR=((-H1*RR+2.0*HH)-((-H1*RP+PH)*RR))/((R*1.0-RR+2.0))
KP=((-H1*RR+2.0*HH)+((-H1*RP+PH)*RR))/((R*1.0-RR+2.0))

EFFECTIVE 1 CALCULATION
YCG=(T0*L*(0-T0/2.0)*L*(0+T0*L/2.0))/((T0*L+0.0+T1[*L])
A=YCG-0.0/2.0
G=C.0
IF(A-G)20,20,21
20 LEFF=(O*O+3.0/12.0+T0*L*(0/2.0-T0/2.0))/(O+1.0-L*(0/2.0)+T1/2.0)
1  *=2.0
G+C 23
21 LEFF=(O*O+3.0/12.0+T0*L*(0/2.0-T0/2.0))/(O+1.0-L*(0/2.0)+T1/2.0)
1 *=2.0
(T0*L+0.0+T1[*L])*(YCG-0/2.0)\*=2.0
23 TS=(T0+T1)/2.0
S=L+A
H=RDReporter
HRM=H/RW
HFS=H/TS
HS=H/S
RTS=3/TS
RETURN
END
1TS,B,P,PF,PPCL,PPC2,PPC3,PPC4,PPC5,PWS,PND)
COMMON NUE,E,T,E2,SIGMAY,LB
REAL NU,NUE,L,NUCA,IEFF,IGMA,L,B,KR,KP
PI=3.1416
P=SIGMAY/(SX+SPHI*SPHI-SX*SPHI)**0.5
PF=SIGMAY/(FX+FX+FPHI*FPHI-FX*FPHI)**0.5
ES=EZ**2.0/ET
FE=FX/SPHI
RM=DM/2.0
TEDA=1.23*(R*T)**C.5/L
NU=0.5-ES*(0.5-NUE)
PPC1=ET*ES**2.0/(12.0*(1.0-NU**2.0)*(SX*PI**2.0-(R**2.0)))
1*(PI**4.0*ET**2.0*ES**2.0*(1.0+C.75*(ES/ET-1.0)))/(0.4*12.0)
2*(1.0+(ES/ET-1.0))**(1.0-NU**2.0)***(1.0+7.5*(ES/ET-1.0))**(-1.0)
PPC2=PI**2.0*ET**2.0/3.0**SX*TEDA**2.0**ET/4.0*(1.0-NU**2.0)**(-0.5)
1=L**2*(1.0+0.75*TEDA*(ES/ET-1.0))/(3.0-2.0**TEDA*(1.0-ET)**(-1.0))
LAMDA=PI*R/LB
DO 7 K=1,5
PPC3=5.0*C.0
N=K
PPC3(K)=2.0*TS*(ET**ES)**0.5*E/RM**((LAMDA**4.0/((N+2.0)
1.0+LAMDA**2.0/2.0)**(N+2.0+LAMDA**2.0)**(1.0-0.75**((
2.0-LAMDA**4.0)**(N+2.0+LAMDA**2.0)**(1.0+2.0)))+
3*(N+2.0-1.0)**ET**EFF/(RM**3.0*L)
C3=PPC3(K)
IF(C3-PPC3)3,3,7
3 PPC3=C3
7 CONTINUE
PPC4=2.0**ET**EFF/(DM**3.0*L*(1.0-NUE**2.0)**ET
PPC5=PI**2.0*ET**B**2.0/(12.0*(1.0-NU**2.0)*(KP+4.0*KR)
1*PI**2.0)
PWS=SIGMAY/(KR+KP-KP-KR)**0.5
PND=4.0*TS**SIGMAY/DM
RETURN
END
STRESS ANALYSIS (STAS) OF WEB STIFFENED SANDWICH

CYLINDRICAL SHELLS
REAL NUE, P, N, MXUO, MPHIOI, MXO1, MPHIOI, RQ1.0, MPHIL10,
MXII, MPHIII, IEFF, KR, KP
DIMENSION RRO(I), RRI(I), TTO(I), TTI(I), BB(I), WL
1 (10), PNUE(10)
READ(5, 70C) KHC
7CC FORMAT(F10.2)
   READ(5, 10C) (RRO(I), RRI(I), TTO(I), TTI(I), BB(I), WL(I),
   1 PNUE(I), 1=1,9)
1CC FORMAT(7F10.2/7F10.2/7F10.2/7F10.2/7F10.2/7F10.2/
   1 7F10.2/7F10.2/7F10.2)
DO 10 I=1,9
   RO=RRO(I)
   RI=RRI(I)
   TC=TTO(I)
   TI=TTI(I)
   B=BB(I)
   L=WL(I)
   NUE=PNUE(I)
WRITE(6, 20C) RO, RI, TO, TI, B, L
200 FORMAT(' STRESS ANALYSIS OF A WEB STIFFENED SANDWICH
   1 CYLINDRICAL PRESSURE HULL'/' OUTSIDE RADIUS (IN.)=','
   2 F10.2' INSIDE RADIUS (IN.)='F10.2' OUTSIDE THICKNESS
   3 (IN)=''F10.2' INSIDE THICKNESS='F10.2' WEB
   4 THICKNESS (IN.)='F10.2' WEB SPACING (IN.)='F10.2'/2x,
   5 ' MXUO' ,2x, ' MPHIOI' ,2x, ' MXO1' ,2x, ' MPHIOI' ,2x,
   6 ' MXII' ,2x, ' MPHIII' ,2x, ' MXII' ,2x, ' MPHIII/'
POP=1.0
THEDAO=(3.0*(1.0-NUE)*2.0)*0.25*L/(RO*TO)*0.50
THEDAI=(3.0*(1.0-NUE)*2.0)*0.25*L/(RI*TI)*0.50
FIJ=THEDAO/2.0*(SINH(THEDAO)+SIN(THEDAO))/(COSH(THEDAO)
1-COST(THEDAO))
FII=THEDAI/2.0*(SINH(THEDAI)+SIN(THEDAI))/(COSH(THEDAI)
1-COST(THEDAI))
8EDAJ=THEDAC/2.0
\[ \text{BEDA1=THEDAI/2.0} \]
\[ \text{E20=THEDAI*(COSH(BEDA1)*SIN(BEDA1)+SINH(BEDA1)*COS(BEDA1))} \]
\[ \text{F21=THEDAI*(COSH(BEDA1)-COS(BEDA1))} \]
\[ \text{F31=0.0*THEDAI/(3.0*(1.0-NUE*2.0))*0.5*(COSH(BEDA1))} \]
\[ \text{1*SIN(BEDAC)-} \]
\[ \text{2*SINN(BEDAI)*COS(BEDAC)/(COSH(BEDAI)-COS(BEDAI))} \]
\[ \text{F41=0.0*THEDAI/(3.0*(1.0-NUE*2.0))*0.5*(COSH(BEDA1))} \]
\[ \text{1*SIN(BEDAI)-} \]
\[ \text{2*SINH(BEDAI)*COS(BEDAI)/(COSH(BEDAI)-COS(BEDAI))} \]
\[ \text{F51=3.0*THEDAI/(3.0*(1.0-NUE*2.0))*0.5*(SINH(BEDAI))} \]
\[ \text{1-SIN(THEDAI)/(COSH(THEDAI)-COS(THEDAI))} \]
\[ \text{F61=3.0*THEDAI/(3.0*(1.0-NUE*2.0))*0.5*(SINH(THEDAI))} \]
\[ \text{GJ=-2.0*R1*2.0*FLI/(TI*4L)} \]
\[ \text{G1=-2.0*R1*2.0*FLI/(TI*VL)} \]
\[ \text{R2=(R0+R1)/2.0} \]
\[ \text{P=(R3-R1)+(TC+TI)/2.0} \]
\[ \text{N=NS[H*(2.0-R1)/(2.0*R1)*60+((2.0*R1)-2.0)*NUE]} \]
\[ \text{L=NS[H*0.0/(2.0-R1)/(2.0*R1)]} \]
\[ \text{P=NS[H*(2.0-R1)+(2.0*R1)-2.0]} \]
\[ \text{J=NS[H*(1.0-R1)/(2.0*R1)]} \]
\[ \text{L=NS[0.0/(2.0-R1)/(2.0*R1)]} \]
\[ \text{UL=(TI/TJ)*VL} \]
\[ \text{H=((G1-0)*(-(R1)*2.0/T0)+(1.0-(NUE/2.0)/(1.0+T1/T0))} \]
\[ \text{1*(R1/R0)))*2.0*N/2.0]*(-N)*((NUE/R0)*2.0/(2.0*T1))} \]
\[ \text{2*(1.0-(1.0)/(1.0+T1/T0)*((R1/R0)))+3.0P/2.0)]/(G1-2)) \]
\[ \text{3 G0-0=}(-N)*(-P)} \]
\[ \text{H1=(T-P)*((R1)*2.0/T0)*((1.0-(NUE/2.0)/(1.0+T1/T0))} \]
\[ \text{1*(R1/R0)))*2.0*N/2.0]-((G0-M)*((NUE/R0)*2.0/(2.0*T1))} \]
\[ \text{2*(1.0-(1.0)/(1.0+T1/T0)*((R1/R0)))+3.0P/2.0])/(-(N)} \]
\[ \text{3 (-P)-((GI-0)*)((GU-M))} \]
\[ \text{SIGNJ=-P*K/R/L/T0} \]
\[ \text{SIGNJ=-P*K/R/L/TI} \]
\[ \text{RL=2.0*E20-NUE*F30} \]
$B11 = 2.0 \times F21 - NUE \times F31$

$B20 = 2.0 \times F20 + NUE \times F30$

$B21 = 2.0 \times F21 + NUE \times F31$

$B3U = F3U$

$B3I = F3I$

$B4C = F4C$

$B50 = 2.0 \times F10 + NUE \times F40$

$B60 = 2.0 \times F10 - NUE \times F40$

$B41 = \Phi 41$

$B51 = 2.0 \times F11 + NUE \times F41$

$B61 = 2.0 \times F11 - NUE \times F41$

C STRESS STATE ONE

$FX00 = (V0/R0-H0) \times B40/L \times SIGU1$

$FPHI0 = (1.0-HC \times B50/L) \times SIGU0$

$MX00 = (V0/R0+H0 \times B50/L) \times SIGU0$

$MPHI0 = (1.0-HC \times B10/L) \times SIGU1$

C STRESS STATE TWO

$FX01 = (V0/R0+H0 \times B40/L) \times SIGU0$

$FPHI1 = (1.0-HC \times B60/L) \times SIGU1$

$MX01 = (V0/R0+H0 \times B20/L) \times SIGU1$

$MPHI1 = (1.0-HC \times B20/L) \times SIGU1$

C STRESS STATE THREE

$FX02 = (V0/R0+H0 \times B40/L) \times SIGU1$

$FPHI2 = (-1.0+H1 \times B41/L) \times SIGU1$

$MX02 = (V0/R0+H0 \times B31/L) \times SIGU1$

$MPHI2 = (-1.0+H1 \times B31/L) \times SIGU1$

C STRESS STATE FOUR

$FX03 = (V0/R0+H0 \times B40/L) \times SIGU1$

$FPHI3 = (-1.0+H1 \times B61/L) \times SIGU1$

$MX03 = (V0/R0+H0 \times B21/L) \times SIGU1$

$MPHI3 = (-1.0+H1 \times B21/L) \times SIGU1$

C WEB STRESS CALCULATION

$H1 = HC + H/2.0$

$PH1 = HI$

$DI = R1 - H1/2.0$

$CR1 = R3 + HC/2.0$
\[
RR = DRI / CRC \\
KR = ((-HH1 x RR x 2.0) + HH0) - ((-HH1 x RR + HH0 x RR) x RR) / (R x (1.0 - RR \\
1 x 2.0)) \\
KP = ((-HH1 x RR x 2.0 + HH1) + ((-HH1 x RR + HH0 x RR) x RR) / (R x (1.0 - RR \\
1 x 2.0)))
\]

**Effective I Calculation**

\[
YCG = (TC x L x (D-T)/2.0) + D x 2.0 x R x 2.0 + T1 x 2.0 x L x 2.0) / (TC x L \\
1 + D x B + T1 x L)
\]

\[
A = YCG - 0.7 x C \\
C = 0.7 \\
10 IF = A - G 20, 20, 21 \\
20 IF = B x D x 3 x C / 12.0 + T1 x L x (D / 2.0 - 10 / 2.0) x 2.0 + T1 x 1 x (D / \\
1 + 2.0 - T1 / 2.0) x 2.0 \\
30 GCL 10 23 \\
21 IF = B x D x 3 x C / 12.0 + T1 x L x (D / 2.0 - T / 2.0) x 2.0 + T1 x 1 x (D / \\
1 + 2.0 - T1 / 2.0) x 2.0 \\
30 WRITE(6, 300) MXC1, */PH100, MX01, */PH101, MX10, */PH110, MX11, \\
1 MPHI11
30 FORMAT(1X, F, 7.2, 1X, F, 7.2, 1X, F, 7.2, 1X, F, 7.2, 1X, F, 7.2, 1X, \\
1F7.2, 1X, F, 7.2, 1X, F, 7.2) \\
40 WRITE(6, 400)
40 FORMAT(2X, 'FX01', 2X, 'FX01', 2X, 'FX01', 2X, 'FX01', 2X, 'FX01', 2X, \\
1, 2X, 'FX10', 2X, 'FX10', 2X, 'FX10', 2X, 'FX10', 2X, 'FX10', 2X, \\
1, 2X, 'FX1', 2X, 'FX1', 2X, 'FX1', 2X, 'FX1', 2X, 'FX1') \\
WRITE(6, 300) FXC1, */PH10C, FXC1, */PH1C1, FXC1, */PH1C1, FXC1, */PH1C1, FXC1, \\
1 EPHI11
15 S = (TC + T1) / 2.0 \\
S = L + B \\
H = R0 - RI \\
P = H / RW \\
PTS = H / TS \\
HS = H / S \\
BTS = B / TS \\
WB = 2.0 + RW / 64.4 x (TC / RG + (RI / R) x (T1 / R0) + (R0 / R) x (B / S)) x \\
1 (D / S)) \\
WRITE(6, 500)
50 FORMAT(5X, 'FX01', 5X, 'FX01', 5X, 'HS', 5X, 'HS', 5X, 'BTS', 5X, 'BTS', 5X,
1  IEFF, 5X, ' 0' />
   WRITE (6, 600) HRW, HS, HTS, B1S, IEFF, 0
6CC FORMAT (2X, 4F10.3, F10.1, F10.3)
   WRITE (6, 800) W
8CC FORMAT (5X, ' W=' F10.3/)
   WRITE (5, 900) KR, KP
9CC FORMAT (5X, ' KR=' F10.3, 5X, ' KP=' F10.3/)
1C CONTINUE
   CALL EXIT
   END
FAILURE MODES CRITERIA (FMCS) FOR FOUR STATES OF STRESS

REAL NU, NUE, L, LAMDA, IEFF, IGMAI, LB, KR, KP
DIMENSION IGMAI(10), EET(10), EES(10), SX(2), SPHI(2),
1 FX(2), FPHI(2), R(2), T(2), VPC(5)
READ(5, 102) (IGMAI(J), EET(J), EES(J), J = 1, 10)
102 FORMAT(3F10.2/3F10.2/3F10.2/3F10.2/3F10.2/3F10.2/3F10.2)
10 READ(5, 100) (SX(I), SPHI(I), FX(I), FPHI(I), R(I), T(I),
1 I = 1, 2)
100 FORMAT(2F10.2/2F10.2/2F10.2/2F10.2/2F10.2/2F10.2)
IF (SX(1) - 1.0) 8, 8, 9
8 CALL EXIT
9 READ(5, 101) IEFF, D, L, NUE, LB, SIGMAY, E, KR, KP, B
101 FORMAT(5F10.2, F10.0, E10.2/3F10.2)
WRITE(6, 103) IEFF, D, L, NUE, LB, SIGMAY, E, B
103 FORMAT(5F10.2, F10.0, E10.2/3F10.2)
WRITE(6, 200)
200 FORMAT(20X, *GRAPHICAL DATA FOR FAILURE MODES'//7X,
1 ' P', ' 7X', ' PF', ' 7X', ' PPC1', ' 6X', ' PPC2', ' 6X', ' PPC3', ' 6X',
2 ' PPC4', ' 6X', ' PPC5', ' 6X', ' SIGMAI'/)
D = R(I) + R(2)
RW = D / 2.0
H = R(I) - R(2)
PI = 3.1417
5 DO 1 I = 1, 2
WRITE(6, 400) I
400 FORMAT(*STATE OF STRESS = 'I2'/)
DO 6 J = 1, 10
SIGMAI = IGMAI(J) * SIGMAY
ET = EET(J)
ES = EES(J) ** 2.0 / ET
P = SIGMAI / (SX(I) * SX(I) + SPHI(I) * SPHI(I) - SX(I) * SPHI(I))
1 ** 0.5
PF = SIGMAI / (FX(I) * FX(I) + FPHI(I) * FPHI(I) - FX(I) * FPHI(I))
1 ** 0.5
NU = 0.5 - ES * (0.5 - NUE)
PPC1 = (ET*EL*2.0/(12.0*(1.0-NU**2.0))*SX(I)*PI**2.0*R
1 *(I)**2.0))
2 *(PI**4.0*R(I)**2.0*T(I)**2.0*(1.0+.75*(ES/ET-1.0)))/
3 *(1.0+(ES/ET-1.0))*(1.0-NU**2.0)/(1.0+.75*(ES/ET-1.0))
FE = SX(I)/SPH(I)
TEDA = 1.23*(R(I)**2.5/L
PPC2 = PI**2.0*E*FE*T(I)/(3.0*SX(I)*TEDA*(1.0-NU**2.0)
1 *(R(I)**2)
2 *(R(I)**2.0*(1.0-0.5/L)**2.0*(1.0+.75*TEDA*(ES/ET-1.0)))/
3 *(3.0-2.0*TEDA*(1.0-FE))
LAMDA = PI*R(I)/LB
DO 7 K = 1,5
PPC3 = 50000.0
N = K
VPC3(K) = (T(I)+T(2))*ET*ES)**0.5*E/RW*(LAMDA**4.0/
1 *(N**2.0-1.0)**
2 LAMDA**2.0/2.0)*(N**2.0+LAMDA**2.0)*(1.0-.75*(1.0-
3 ES/ET)*LAMDA**2.
4 0.0/(N**2.0+LAMDA**2.0)**2.0))+(N**2.0-1.0)*ET*E*
5 IEFF/(RW**3.0*L)
C3 = VPC3(K)
IF(C3-PPC3)3,3,7
3 PPC3 = C3
7 CONTINUE
PPC4 = 4.0*E*IEFF/(DM**3.0*L*(1.0-NU**2.0))*ET
PPC5 = (4.0*(PI**2.0)*E*ET*B**2.0)/(12.0*(1.0-NU**2.0)*
1 *(KP**4.0*KR)**2.0)
PND = 2.0*(T(I)+T(2))*SIGMAY/DM
PWS = SIGMAY/(KR*KR+KP*KP-KR*KP)**0.5
WRITE(6,500) P, PF, PPC1, PPC2, PPC3, PPC4, PPC5, SIGMAI
5CC FORMAT(2X,F10.1,1X,F10.1,1X,F10.1,1X,F10.1,1X,F10.1,1X,
1 F10.1,1X,F10.1,1X,F10.2/)
6 CONTINUE
5 CONTINUE
WRITE(6,501) PWS, PND
x in o < o => O O - O LL O U - o
RESULTS

1. Geometric similitude was determined to exist. The stress analysis and failure mode programs output are in Appendix D for three geometrically similar structures. Graph 1 shows the failure modes for these three geometrically similar structures made from HY-160 steel. The effect of the inelastic range of this material on the buckling modes is indicated. The example has underdeveloped webs (PWS, PPC4).

2. Maximum stress states were found to be located as shown below.

3. The results are plotted in graphical form. The results are divided basically to show (a) the failure modes of the shells, (b) the failure modes of the webs, and (c) optimum proportions and weight/displacement results. The presentation of the results in this form was (1) to facilitate comparison with ring-stiffened cylinders, and (2) to
show the relationship of the various failure modes. The form of the results are very similar to those of ring-
stiffened cylinders.

4. Graphs 2 through 7 are sample results of the stress and buckling analysis for the shell portion of the web-
stiffened sandwich cylinder using HY-30 steel.

5. Graphs 8 through 11 are sample results for the web portion non-dimensionalized by $P_{ND}$ where

$$P_{ND} = \frac{2(T_0 + T_1) \sigma_y}{D_M}$$

The material was HY-80 steel.

6. Graphs 12 and 13 show the effect of yield strength and the modulus of elasticity on shell failure modes respectively.

7. Graphs 14 and 15 and Tables 1 and 2 are concerned with the optimization results. The optimum proportions for ring-
stiffened cylinders shown in Graph 15 were based upon using 92, 92A criteria for shell yield (outer fibers reaching yield stress).

8. For a given hydrostatic loading and hull diameter the shell thicknesses for the web-stiffened sandwich are approximately half that of a ring-stiffened cylinder.
FAILURE MODES
INELASTIC ANALYSIS
HY-160 STEEL

Graph 1

Collapse Pressure "p" pounds per sq. in.

State of Stress (σ2)

PCPC5
PCPC3
PCPC1
PCPC2
PCPC4

PND
PF
PWS
Graph 2

Collapse Pressure "P" Pounds per sq. in.

Axisymmetric Mode
\[ \frac{h}{R_w} = 0.10 \]
\[ \frac{b}{D_m} = 0.0025 \] (PCC1)

Unsupported Length ÷ Diameter \( \left( \frac{L}{D_m} \right) \)
GRAPH 3

COLLAPSE PRESSURE \( \frac{P}{\text{POUNDS PER SQ. IN.}} \)

\[ \frac{t_o}{D_m} = 0.0100 \]
\[ \frac{t_o}{D_m} = 0.0050 \]
\[ \frac{t_o}{D_m} = 0.0025 \]

ASYMMETRIC MODE

\[ \frac{h}{R_w} = 0.10 \]
\[ \frac{b}{D_m} = 0.0025 \]

(PPC 2)
Graph 4

Shell Yield

\[ h/R_w = 0.10 \]
\[ b/D_m = 0.0025 \]

Collapse Pressure "P" Pounds per sq. in.

Unsupported Length ÷ Diameter \( (L/D_m) \)
Composit of Shell
Failure Modes

\[ \frac{h}{R_w} = 0.10 \]
\[ \frac{b}{D_m} = 0.0025 \]

Graph 5

Collapse Pressure \( P \), Pounds Per Sq. In.

Unsupported Length ÷ Diameter \( \left( \frac{L}{D_m} \right) \)
GRAPH 6

COLLAPSE PRESSURE $P$ POUNDS PER SQ. IN.

UNSUPPORTED LENGTH $L$ / DIAMETER $D_m$

SHELL YIELD, $P$

- $h/R_w = 0.10$
- $b/D_m = 0.0025$

- dash: inner
- solid: outer

$t/D_m = 0.00125$
$t/D_m = 0.0025$
$t/D_m = 0.0050$
$t/D_m = 0.0100$
Graph 7

Collapse pressure "P" (in pounds per sq. in.)

- For $b/t_s = 0.500$, $h/R_w = 0.020$
- For $b/t_s = 0.786$, $h/R_w = 0.020$

- For $b/t_s = 0.500$, $h/R_w = 0.058$
- For $b/t_s = 0.786$, $h/R_w = 0.058$

Shell Yield, P

Effect of $b$, $h$, (webs) on cross-over of inner and outer shell stress

Unsupported Length ÷ Diameter ($L/D_m$)
Graph 10

EFFECT OF WEBS
\[ \frac{t_a}{D_m} = \frac{t_i}{D_m} = 0.010 \]
\[ \frac{h}{t_s} = 5.00 \]
\[ \frac{L}{D_m} = 0.10 \]

WEB THICKNESS / SHELL AVERAGE THICKNESS, \( \frac{b}{t_s} \)
COMPOSITE OF WEB FAILURE MODES

- $t_0/d_m = t_i/d_m = 0.005$
- $h/t_s = 5.00$
- $b/t_s = 0.50$

UNLIMITED LENGTH ÷ DIAMETER ($L/D_m$)
EFFECT OF STRENGTH ($G_Y$)
COLLAPSE PRESSURE "P"  POUNDS PER SQ. IN.

UNSUPPORTED LENGTH / DIAMETER (L/D) vs. DIAMETER (D)

GRAPH 13

EFFECT OF ELASTICITY (E)
NON-DIMENSIONAL WEB DEPTH

(CR/Rw)

OPTIMUM WEB DEPTH (h)

O---O HY-80 STEEL

A---A HY-160 STEEL

COLLAPSE PRESSURE "P" POUNDS PER SQ. IN.

GRAPH 14
Graph 15

Weight/Displacement vs. Collapse Depth (Pressure)

- O - O HY-80 Steel
- △ - △ HY-160 Steel

Collapse Pressure "p" Pounds per Sq. In.

Weight/Displacement (W/D)

Unstiffened Cylinder

Ring Stiffened Cylinder HY-80
<table>
<thead>
<tr>
<th>$P_{coll}$</th>
<th>$h/R_w$</th>
<th>$h/s$</th>
<th>$h/t_s$</th>
<th>$b/t_s$</th>
<th>$t_o/D$</th>
<th>$t_i/D$</th>
<th>$L/D$</th>
<th>$W/D$</th>
</tr>
</thead>
<tbody>
<tr>
<td>500</td>
<td>0.058</td>
<td>1.43</td>
<td>17.65</td>
<td>0.588</td>
<td>0.0019</td>
<td>0.0015</td>
<td>0.020</td>
<td>0.137*</td>
</tr>
<tr>
<td>1000</td>
<td>0.095</td>
<td>1.56</td>
<td>17.39</td>
<td>0.696</td>
<td>0.003</td>
<td>0.00275</td>
<td>0.030</td>
<td>0.239</td>
</tr>
<tr>
<td>2000</td>
<td>0.113</td>
<td>2.14</td>
<td>12.00</td>
<td>0.600</td>
<td>0.005</td>
<td>0.005</td>
<td>0.025</td>
<td>0.436</td>
</tr>
<tr>
<td>5000</td>
<td>0.148</td>
<td>2.16</td>
<td>6.15</td>
<td>0.538</td>
<td>0.014</td>
<td>0.014</td>
<td>0.030</td>
<td>1.061</td>
</tr>
</tbody>
</table>

### HY-150

<table>
<thead>
<tr>
<th>$P_{coll}$</th>
<th>$h/R_w$</th>
<th>$h/s$</th>
<th>$h/t_s$</th>
<th>$b/t_s$</th>
<th>$t_o/D$</th>
<th>$t_i/D$</th>
<th>$L/D$</th>
<th>$W/D$</th>
</tr>
</thead>
<tbody>
<tr>
<td>500</td>
<td>0.077</td>
<td>2.50</td>
<td>33.02</td>
<td>0.976</td>
<td>0.011</td>
<td>0.010</td>
<td>0.015</td>
<td>0.125</td>
</tr>
<tr>
<td>1000</td>
<td>0.077</td>
<td>1.87</td>
<td>26.67</td>
<td>0.933</td>
<td>0.016</td>
<td>0.014</td>
<td>0.020</td>
<td>0.155</td>
</tr>
<tr>
<td>2000</td>
<td>0.077</td>
<td>1.73</td>
<td>14.29</td>
<td>0.893</td>
<td>0.028</td>
<td>0.028</td>
<td>0.020</td>
<td>0.258</td>
</tr>
<tr>
<td>5000</td>
<td>0.113</td>
<td>2.14</td>
<td>9.23</td>
<td>0.462</td>
<td>0.008</td>
<td>0.005</td>
<td>0.025</td>
<td>0.522</td>
</tr>
</tbody>
</table>

*This value seemed high and thus a scatter check was conducted about this point. Results of this check are in Table 2.*
Table 2

OPTIMUM SCATTER ABOUT

$P_{\text{coll}} = 500 \text{ psi}$

<table>
<thead>
<tr>
<th>$P_{\text{coll}}$</th>
<th>Mode</th>
<th>$h/R_N$</th>
<th>$t_o/b$</th>
<th>$t_i/D$</th>
<th>$L/D$</th>
<th>$b/t_S$</th>
<th>W/D</th>
<th>W/D*</th>
</tr>
</thead>
<tbody>
<tr>
<td>508</td>
<td>PF</td>
<td>.058</td>
<td>.0014</td>
<td>.0014</td>
<td>.025</td>
<td>.786</td>
<td>.124</td>
<td>.122</td>
</tr>
<tr>
<td>554</td>
<td>PWS</td>
<td>.058</td>
<td>.0014</td>
<td>.0014</td>
<td>.020</td>
<td>.643</td>
<td>.128</td>
<td>.115</td>
</tr>
<tr>
<td>524</td>
<td>PWS</td>
<td>.058</td>
<td>.0015</td>
<td>.0014</td>
<td>.025</td>
<td>.759</td>
<td>.127</td>
<td>.121</td>
</tr>
<tr>
<td>581</td>
<td>PWS</td>
<td>.058</td>
<td>.0016</td>
<td>.0014</td>
<td>.020</td>
<td>.600</td>
<td>.134</td>
<td>.115</td>
</tr>
<tr>
<td>647</td>
<td>PPC4</td>
<td>.058</td>
<td>.0018</td>
<td>.0014</td>
<td>.020</td>
<td>.625</td>
<td>.145</td>
<td>.112</td>
</tr>
<tr>
<td>677</td>
<td>PPC4</td>
<td>.058</td>
<td>.0018</td>
<td>.0015</td>
<td>.020</td>
<td>.667</td>
<td>.153</td>
<td>.113</td>
</tr>
</tbody>
</table>

*Normalized W/D by ($\frac{500}{P_{\text{coll}}}$).
DISCUSSION OF RESULTS

1. High strength materials having inelastic zones are more prone to failure by buckling than elastic materials because of the reduction of the value of $E$ above the proportional limit.

2. Graphs 2 through 5 indicate that close spacing (small $L$) of the annular rings (webs) is required in order to utilize the benefits of this type of structure. This follows from the more uniform stressing in this configuration.

3. Maximum stresses in the shell occur in the outer shell ($T_o = T_i$) when web spacing ($L$) is large. However, when the spacing is small the stress in the inner shell can exceed the outer. This occurs when the web thickness and web depth is insufficient to carry the load (see Graphs 6 and 7). Graph 6 suggests that a more evenly distributed stress can be obtained by making $T_o > T_i$ for thin shells. Optimization bears this out (Table 1).

4. Graphs 8 through 11 indicate that web strength and buckling determine web depth and thickness requirements and have a critical effect on the weight/displacement ratio ($PWS = f(bh)$, $PPC4 = f(bh^3)$, $PPC5 = f(b/h)^2$).
5. Shell buckling modes (PPC1, PPC2) become prevalent only in (a) thin shells -- either shallow depth load design or high strength materials, or (b) materials with low moduli of elasticity (E) such as titanium or aluminum, or (c) in the inelastic range of high strength materials.

6. Axisymmetric failure mode (PPC1) is prevalent only in very thin shells \(t/D_m < .0015\).

7. Higher strength materials and materials with low moduli of elasticity should have smaller web spacing in the optimum proportions (PPC2 shifts to the left -- thinner shells for the high strength materials; smaller E for those with low modulus of elasticity).

8. Table 1 and Graph 14 indicate that web depth increases with depth for the optimum scantlings. This is seen from the need to increase the web size with loading. Also \(h/t_s\) decreases with depth indicating \(t_s\) is increasing faster than web depth, \(h\). The parameters \(h/s\) and \(b/t_s\) do not change appreciably.

9. Optimization scatter about the optimum scantlings indicates more than one optimum configuration (Table 2).

10. Graph 15 shows that there is an advantage of this structure over the ring-stiffened cylinder of some order. Further examination is required using membrane yield in
place of yielding at the outer fibers to compare to the optimized results of ring-stiffened cylinders using Lunchick's formulation for shell yield. Also, comparison with experimental failures is required to verify the formulation used in this investigation or to justify the use of membrane stresses.

11. The advantages of thinner shells (homogeneity of material, greater ductility, better notch toughness, and easier to form or weld) compared to those of the ring-stiffened cylinder make the web-core sandwich cylindrical hull very appealing.

12. The results set forth are thought to be conservative. Yield (P, PF, PWS) at the outer fibers was the criteria of yield failure. If membrane stresses were permitted to reach yielding more dramatic results would have ensued.
CONCLUSIONS

1. Web-core sandwich type cylindrical pressure hulls possess structural efficiencies in the order of 10-20% higher than those of conventional ring-stiffened construction at depths of 2,000 feet and 5% higher at depths of 10,000 feet.

2. Higher strengths realized with the sandwich designs over those of conventional ring-stiffened cylinders can be explained:

   a. Symmetry of the sandwich cross section allows a more uniform and more complete stressing of the available material. The outstanding flange of a T-frame on a conventional ring-stiffened pressure hull is stressed only in the circumferential direction, i.e., a uniaxial state of stress exists. In a sandwich pressure hull, the "flange" material is in reality a second shell which is more efficiently stressed in both the axial and circumferential directions, i.e., a biaxial state of stress exists.

2. Sandwich geometry is inherently more stable structurally permitting straining of the material well into the work-hardening range with concomitant beneficial effects of higher strength levels (inelastic materials).
3. Inelastic analysis is very necessary in the use of high strength materials.

4. Sandwich construction permits the use of thinner plating material with its superior homogeneous, higher strength, greater ductility, and better notch toughness characteristics over thicker plating. In addition, thinner plating is easier to form and weld.

5. The use of this type of structure is appealing in several designs:
   (a) Large diameters hulls -- thinner plating.
   (b) Deep depth hulls -- thinner plating.
   (c) Shallow depth hulls (<2000 feet) -- 10-20% more efficient structure.
RECOMMENDATIONS

1. Utilize the membrane stresses developed in Appendix A in the stress analysis and failure mode programs in place of the outer fiber stresses.

2. Compare experimental results to theoretical and determine whether membrane stresses or outer fiber stresses are more realistic.

3. Formulate a more sophisticated optimization program utilizing a search technique (reduce program cost).
References


APPENDIX A

AXISYMMETRIC STRESSES
IN AN ISOTROPIC, WEB-STIFFENED SANDWICH CYLINDER
LOADED WITH UNIFORM EXTERNAL PRESSURE

BASIC EQUATIONS

For cylindrical shells loaded under uniform pressure, the three principal axes of strain are

\[ \varepsilon_x = \frac{1}{E} \left[ \sigma_x - \nu (\sigma \phi + \sigma_r) \right] \]
\[ \varepsilon_\phi = \frac{1}{E} \left[ \sigma_\phi - \nu (\sigma_x + \sigma_r) \right] \]  \hspace{1cm} (1)
\[ \varepsilon_r = \frac{1}{E} \left[ \sigma_r - \nu (\sigma_x + \sigma_\phi) \right] \]

Since the radial stress in the shell is generally small in comparison to the longitudinal and circumferential stresses, Equation (1) reduces to

\[ \varepsilon_x = \frac{1}{E} \left[ \sigma_x - \nu \sigma_\phi \right] \]
\[ \varepsilon_\phi = \frac{1}{E} \left[ \sigma_\phi - \nu \sigma_x \right] \]  \hspace{1cm} (2)

for the longitudinal and circumferential strains, respectively.
The equilibrium of forces and moment of forces in the element requires that

\[ \frac{dQ_x}{dx} = \frac{N_\phi}{R} = -P \]

(3)

\[ \frac{dM_x}{dx} = Q_x \]

if the beam-column effect due to the axial portion of the hydrostatic pressure is neglected.

The moment curvature relationship

\[ M_x = -D \frac{d^2w}{dx^2} \]

where

\[ D = \frac{Et^3}{12(1-\nu^2)} \]

(4)
Equations (3) and (4) can be combined to form the following general differential equation of equilibrium

\[ D \frac{d^4 w}{dx^4} + \frac{N_\phi}{R} = P \]  

(5)

The strains in equation (2) can be expressed in terms of the forces \( N_x \) and \( N_\phi \) as

\[ \varepsilon_x = \frac{1}{E t} (N_x - vN_\phi) \]

\[ \varepsilon_\phi = \frac{1}{E t} (N_\phi - vN_x) \]  

(6)

The circumferential strain can also be expressed as

\[ \varepsilon_\phi = \frac{w}{R} \]

(7)

By substituting Equations (6) and (7) into Equation (5), the general equation of equilibrium is

\[ D \frac{d^4 w}{dx^4} + \frac{E t}{R^2} w = P - \frac{v}{R} N_x \]  

(8)

The edge conditions at \( X = 0 \) and \( X = L \) are

\[ \frac{dw}{dx} \bigg|_{x=0} = \frac{dw}{dx} \bigg|_{x=L} = 0 \]

and

\[ -\frac{d^3 w}{dx^3} \bigg|_{x=0} = \frac{d^3 w}{dx^3} \bigg|_{x=L} = \frac{Q_x}{D} \]  

(9)

With these boundary conditions, the general solution of equation (8) is
where

\[ w = \frac{2R^2Q_x}{EtL} \frac{\theta}{2} \left[ B_1B_3 + B_4 - B_5 - B_2B_6 \right] + \frac{R^2P}{Et} - \frac{\nu N_xR}{E} \frac{t}{t} \]  

\( \theta = \sqrt[4]{\frac{3(1-v^2)}{t^2R^2} L} \)

\[ B_1 = \frac{\sinh \theta + \sin \theta}{\cosh \theta - \cos \theta} \]

\[ B_2 = \frac{\sinh \theta - \sin \theta}{\cosh \theta - \cos \theta} \]

\[ B_3 = \cos \left( \frac{\theta_x}{L} \right) \cosh \left( \frac{\theta_x}{L} \right) \]

\[ B_4 = \sin \left( \frac{\theta_x}{L} \right) \cosh \left( \frac{\theta_x}{L} \right) \]

\[ B_5 = \cos \left( \frac{\theta_x}{L} \right) \sinh \left( \frac{\theta_x}{L} \right) \]

\[ B_6 = \sin \left( \frac{\theta_x}{L} \right) \sinh \left( \frac{\theta_x}{L} \right) \]

**STRESSES AND STRAINS**

The longitudinal stresses on the shell section with externally applied pressure for the outer and inner surfaces are

\[ \sigma_x = \frac{N_x}{t} + \frac{6M_x}{t^2} \]  

\[ \sigma_x = \frac{N_x}{t} + \frac{6M_x}{t^2} \]
From Equations (4) and (10), the longitudinal stresses of Equation (11) are

\[ \sigma_x = \frac{R}{t} \left( \frac{N_x}{R} + \frac{Q_x G_1}{L \psi} \right) \]  \hspace{1cm} (12)

From Equations (2), (7), (10), and (12), the circumferential stresses are

\[ \sigma_\phi = -\frac{R}{t} \left[ -p + \frac{Q_x}{L} \left( G_2 + \nu \frac{G_1}{\psi} \right) \right] \]  \hspace{1cm} (13)

where

\[ G_1 = \theta \left( B_2 B_3 - B_4 - B_5 + B_1 B_6 \right) \]

\[ G_2 = \theta \left( B_1 B_3 + B_4 - B_5 - B_2 B_6 \right) \]  \hspace{1cm} (14)

\[ \psi = \sqrt{\frac{1 - \nu^2}{3}} \]

DEFORMATION OF ISOTROPIC CIRCULAR RINGS

For a web-stiffened sandwich cylinder, the outside and inside shells are separated by a circular ring. A segment for this type of ring is shown below.
Neglecting stresses in the X-direction for the ring, Equation (1) reduces to

\[ \varepsilon_{\phi} = \frac{1}{E} \left( \sigma_{\phi} - \nu \sigma_{r} \right) \]
\[ \varepsilon_{r} = \frac{1}{E} \left( \sigma_{r} - \nu \sigma_{\phi} \right) \]  

(15)

These strains can also be expressed as

\[ \varepsilon_{\phi} = \frac{\hat{W}}{R} \]
\[ \varepsilon_{r} = \frac{d\hat{W}}{dr} \]  

(16)

The equilibrium of radial forces on the small element requires that

\[ \sigma_{\phi} - \sigma_{r} - R \frac{d\sigma_{r}}{dr} = 0 \]  

(17)
With the relationships in Equations (15) and (16), the general solution of Equation (17) is

\[ W(R) = \frac{-Q_i \left( \frac{R_{RI}}{R_{RO}} \right)^2 + Q_o \left( \frac{R}{R_{RO}} \right)}{(1 - \nu)} \]

\[ + \frac{-Q_i \left( \frac{R_{RI}}{R_{RO}} \right) + Q_o \left( \frac{R_{RI}}{R_{RO}} \right)}{(1 - \nu)} \left( \frac{R_{RI}}{R} \right) \]

\[ \frac{Z}{2} \left( \frac{b}{R_{RO}} \right) \left[ 1 - \left( \frac{R_{RI}}{R_{RO}} \right)^2 \right] \]

where

\[ Z = \frac{E}{1 - \nu^2} \]

\[ R_{RO} = R_O + \frac{T_O}{2} \]

\[ R_{RI} = R_I - \frac{T_I}{2} \]

\[ R_{RI} = \frac{R_{RI}}{R_{RO}} \]

By differentiating Equation (18) and by combining Equations (15), (16), and (17), the circumferential stress is

\[ \sigma_\phi(R) = \frac{-Q_i (R_{RI})^2 + Q_o}{b \left( \frac{R}{R_{RO}} \right)} + \frac{-Q_i (R_{RI}) + Q_o (R_{RI})}{2 \left( \frac{R}{R_{RO}} \right) \left[ 1 - (R_{RI})^2 \right]} \left( \frac{R_{RI}}{R} \right) \]

\[ \frac{b}{\left( \frac{R}{R_{RO}} \right) \left[ 1 - (R_{RI})^2 \right]} \]
and the radial stress is

\[
\sigma_r(R) = \frac{\left[-Q_1(RR)^2 + Q_0\right]\left(\frac{R}{RRO}\right) - \left[-Q_1(RR) + Q_0(RR)\right]\left(\frac{RRI}{R}\right)}{b\left(\frac{R}{RRO}\right)\left[1 - (RR)^2\right]} \tag{20}
\]

**DEFORMATIONS OF SANDWICH SHELLS**

Before determining the stresses in the cylindrical shells of a sandwich structure, it is first necessary to obtain the forces \(Q_0\) and \(N_0\) on the outside shell and the forces \(Q_i\) and \(N_i\) on the inside shell.

The radial deflection of the outside shell at the web-shell juncture can be expressed in terms of the forces as

\[
W_{OS} = g_0Q_0 + d_0N_0 + f_0P \tag{21}
\]
and the radial deflection of the inside shell

\[ W_{is} = g_{i1}Q_1 + d_{i1}N_1 \]  

(22)

Similarly, the radial deflection of the outer fiber of the web at the web-shell juncture can be expressed as

\[ W_{ow} = g_{oo}\left( Q_o + \frac{D}{2} P \right) + g_{ol}Q_1 \]  

(23)

and the radial deflection of the inner fiber can be expressed as

\[ W_{iw} = g_{io}\left( Q_o + \frac{D}{2} P \right) + g_{ii}Q_1 \]  

(24)

Since the stresses in a sandwich structure are linear functions of the applied pressure, the coefficients of the force in Equations (21), (22), (23), and (24) will be developed for a unit of applied pressure \((p = 1.0)\). By substituting the geometric properties of the outside shell into Equation (10), the coefficients of the forces in Equation (21) are

\[ g_o = -\frac{2RO^2\theta_o}{2ETOL}\frac{\sinh\delta_o + \sin\theta_o}{\cosh\theta_o - \cos\theta_o} \]  

(25)

\[ d_o = -\frac{vRO}{ETO} \]

\[ f_o = \frac{RO^2}{ETO} \]

where

\[ \theta_o = \sqrt{3(1-v^2)} \frac{L_o}{KD \cdot TO} \]
Also by the geometric properties of the inside shell, the coefficients of the forces in Equation (22) are

\[ g_i = -\frac{2RI^2}{2ETIL} \left( \frac{\sinh \theta_i + \sin \theta_i}{\cosh \theta_i - \cos \theta_i} \right) \]

\[ d_i = -\frac{\nu RI}{ETI} \]

where

\[ \theta_i = \frac{4L}{\sqrt{3(1-\nu^2)}} \]

The hyperbolic and trigonometric function in Equations (25) and (26) are redefined by

\[ F = \frac{\theta}{2} \left[ \frac{\sinh \theta + \sin \theta}{\cosh \theta - \cos \theta} \right] \]

By Equation (27), the coefficients \( g_0 \) and \( g_i \) are

\[ g_0 = -\frac{2R^2}{ETOL} F_0 \]

\[ g_i = -\frac{2RI^2}{ETIL} F_i \] (28)

By substituting combinations of unit and zero forces into Equation (17), the coefficients in Equations (23) and (24) are

\[ g_{oo} = CRRO \left[ (1 + RR^2) - \nu(1 - RR^2) \right] \]

\[ g_{oi} = -2CRRI(RR) \]

\[ g_{io} = 2CRRO(RR) \]

\[ g_{ii} = -CRRI \left[ (1 + RR^2) - \nu(1 - RR^2) \right] \] (29)
where

\[ C = \frac{2}{bE(1 - \frac{R}{R'})} \]

From the equilibrium conditions of the end forces on a sandwich cylinder loaded under an external uniform pressure, the relationship between the outside and inside axial forces for \( p = 1.0 \) psi is

\[ N_o R_o + N_i R_i = \frac{R O^2}{2} \]  

(30)

By assuming that the axial membrane strain in the outside shell is equal to the axial membrane strain in the inside shell, i.e., \( \frac{N_o}{E T O} = \frac{N_i}{E T I} \), the axial forces are

\[ N_o = \frac{R O}{2} \left( 1 + \frac{T I \cdot R I}{T O \cdot R O} \right) \]  

(31)

\[ N_i = \frac{T I}{T O} N_o \]

For compatibility of deformations in the radial direction, the relationships that must exist are \( W_{os} = W_{ow} \) and \( W_{is} = W_{iw} \). With these relationships and Equations (21) to (29), the radial forces \( Q_o \) and \( Q_i \) can be expressed in terms of the two following simultaneous equations:
\[(g_O - g_{oo}) Q_o - g_{oi} Q_i = -\frac{RO^2}{ETO} \left[ 1 - \frac{v/2}{1 + \frac{TI}{RO}} \right] \]

\[\frac{b}{2} g_{oo} \]

\[ -g_{io} Q_o + (g_i - g_{ii}) Q_i = \frac{v}{2E} \frac{RO^2}{TI} \left[ 1 - \frac{1}{1 + \frac{TI}{RO}} \right] \]

\[\frac{b}{2} g_{io} \]

(32)

The forces \(N_o, N_i, Q_o,\) and \(Q_i\) can be computed from Equations (31) and (32).

**STRESS DISTRIBUTION IN SANDWICH STRUCTURES**

The stresses on the outside shell of the isotropic sandwich are determined by substituting \(N_o, Q_o,\) and the geometric properties of the outside shell into Equations (12) and (13). From these relationships, the circumferential stress in the outside shell on an outer fiber is

\[\sigma_{\phi oo} = -P \frac{RO}{TO} \left[ -1 + \frac{Q_o}{L} \left( G_{20} + \frac{v}{\psi} G_{10} \right) \right] \]

on an inner fiber is

\[\sigma_{\phi oi} = -P \frac{RO}{TO} \left[ -1 + \frac{Q_o}{L} \left( G_{20} - \frac{v}{\psi} G_{10} \right) \right] \]

and the membrane stress is

\[\sigma_{\phi om} = -P \frac{RO}{TO} \left[ -1 + \frac{Q_o}{L} G_{20} \right] \]

(33)
The **longitudinal stress in the outside shell** on an outer fiber is

\[ \sigma_{xoo} = P \frac{R_O}{T_O} \left[ \frac{N_O}{R_O} - \frac{Q_O}{L\psi} G_{10} \right] \]

on an inner fiber

\[ \sigma_{xoi} = P \frac{R_O}{T_O} \left[ \frac{N_O}{R_O} + \frac{Q_O}{L\psi} G_{10} \right] \]

and the membrane stress is

\[ \sigma_{xom} = P \frac{N_O}{T_O} \]

And likewise the stresses on the inside shell are determined by substituting \( N_i, Q_i, \) and the geometric properties of the inside shell into Equations (12) and (13).

\[ \sigma_{\phi io} = -P \frac{R_I}{T_I} \frac{Q_i}{L} \left[ G_{2i} + \frac{\nu}{\psi} G_{li} \right] \]

\[ \sigma_{\phi ii} = -P \frac{R_I}{T_I} \frac{Q_i}{L} \left[ G_{2i} - \frac{\nu}{\psi} G_{li} \right] \]

\[ \sigma_{\phi im} = -P \frac{R_I}{T_I} \frac{Q_i}{L} G_{2i} \]

\[ \sigma_{xio} = P \frac{R_I}{T_I} \left[ \frac{N_i}{R_I} - \frac{Q_i}{L\psi} G_{li} \right] \]

\[ \sigma_{xii} = P \frac{R_I}{T_I} \left[ \frac{N_i}{R_I} + \frac{Q_i}{L\psi} G_{li} \right] \]

\[ \sigma_{xim} = P \frac{N_i}{T_I} \]

(34)
The stresses in the web stiffeners at the outside and inside shell locations are determined by substituting the outside radius RRO and inside radius RRI into Equations (19) and (20). Also, the total force on the outer surface of the web consists of the radial force $Q_o$ at the web-shell juncture and the force attributed to the applied pressure directly on the web. The total forces ($H_o$ and $H_i$) on the outer and inner surfaces of the web for a unit pressure ($p = 1.0$ psi) are

\[
H_o = Q_o + b/2
\]

\[
H_i = Q_i
\]

With these relationships, the circumferential stress in the web on the outer surface is

\[
\sigma_{\phi w_0} = -p \frac{\left[(-H_i(RR)^2 + H_o) + (-H_i(RR) + H_o(RR))RR\right]}{b/2 (1 - RR^2)}
\]

and the radial stress

\[
\sigma_{r_0 w_0} = -p \frac{\left(-H_i(RR)^2 + H_o\right) - \left(-H_i(RR) + H_o(RR)\right)RR}{b/2 (1 - RR^2)}
\]
and on the inner surface

\[
\sigma_{\phi wi} = \frac{\left[ (-H_1(RR)^2 + H_o)_{(RR)} - (-H_1(RR) + H_o(RR)) \right]}{\frac{b}{2} (RR) (1-RR^2)}
\]

\[
\sigma_{rwi} = \frac{\left[ (-H_1(RR)^2 + H_o)_{(RR)} - (-H_1(RR) + H_o(RR)) \right]}{\frac{b}{2} (RR) (1-RR^2)}
\]

(39)
APPENDIX B

BUCKLING ANALYSIS

1. Inelastic Asymmetric (Lobar) Buckling Analysis.

The buckling equations for a fully plastic cylinder specialized for the case of hydrostatic pressure loading

\[
\frac{E_T}{2E_s}\left(\frac{1}{2} \frac{\partial^2 u}{\partial x^2} + \frac{1}{R} \frac{\partial^2 v}{\partial x \partial \phi} + \frac{1}{R} \frac{\partial w}{\partial x}\right)
\]

\[
+ \frac{3}{4} \frac{\partial^2 u}{\partial x^2} + \frac{1}{4R^2} \frac{\partial^2 u}{\partial \phi^2} + \frac{1}{4R} \frac{\partial^2 v}{\partial x \partial \phi} = 0
\]

\[
\frac{E_T}{E_s} \left(\frac{1}{2} \frac{\partial^2 u}{\partial x \partial \phi} + \frac{1}{R^2} \frac{\partial^2 v}{\partial \phi^2} + \frac{1}{R} \frac{\partial w}{\partial \phi}\right)
\]

\[
+ \frac{1}{4} \frac{\partial^2 v}{\partial x^2} + \frac{1}{4R} \frac{\partial^2 u}{\partial x \partial \phi} = 0
\]

\[
\frac{4}{3} \frac{E_s T}{R} \left(\frac{E_T}{E_s}\right) \left(\frac{1}{2} \frac{\partial u}{\partial x} + \frac{1}{R} \frac{\partial v}{\partial \phi} + \frac{w}{R}\right)
\]

\[
+ D \left[\frac{E_T}{E_s} \left(\frac{1}{4} \frac{\partial^4 u}{\partial x^4} + \frac{1}{R^2} \frac{\partial^4 w}{\partial x^2 \partial \phi^2} + \frac{1}{R^4} \frac{\partial^4 w}{\partial \phi^4}\right)
\]

\[
+ \frac{3}{4} \frac{\partial^4 w}{\partial x^4} + \frac{1}{R^2} \frac{\partial^4 w}{\partial x^2 \partial \phi^2}\right] + N_x \frac{\partial^2 w}{\partial x^2}
\]

\[
+ N_\phi \frac{1}{R^2} \frac{\partial^2 w}{\partial \phi^2} + P = 0
\]
where \( x \) and \( \phi \) are respectively the axial and circumferential coordinates,

\[ u, v, \text{ and } w \] are the axial, tangential, and radial displacements,

\( E_s \) and \( E_T \) are the secant and tangent moduli,

\( R \) is the radius to shell mid-surface,

\( t \) is the shell thickness,

\( v \) is Poisson's ratio,

\( D \) is the bending rigidity = \( E_s t^3/12(1-v^2) \)

\( N_x \) and \( N_\phi \) are forces per unit length in the axial and circumferential directions, and

\( p \) is the hydrostatic pressure.

With several differentiations, \( u \) and \( v \) can be eliminated from Equations (1a), (1b) and (1c) so that a single eight-order equation in \( w \) is obtained:

\[
D \left[ \frac{E_T}{E_s} \nabla^8 w + \left( 1 - \frac{E_T}{E_s} \right) \left( \nabla^4 \left( \frac{3}{2} \frac{\partial^4 w}{\partial x^4} + \frac{1}{R^2} \frac{\partial^4 w}{\partial x^2 \partial \phi^2} \right) \right) \right]
\]

\[ + \frac{3}{4} \left( \frac{E_s}{E_T} - 1 \right) \left( \frac{3}{4} \frac{\partial^8 w}{\partial x^6} + \frac{1}{R^2} \frac{\partial^8 w}{\partial x^2 \partial \phi^2} \right) \]

\[ + \frac{E_s t}{R^2} \frac{\partial^4 w}{\partial x^4} + N_x \left[ \nabla^4 \frac{\partial^2 w}{\partial x^2} + \frac{3}{4} \left( \frac{E_s}{E_T} - 1 \right) \frac{\partial^6 w}{\partial x^6} \right] \]

\[ + N_\phi \left[ \nabla^4 \frac{1}{R^2} \frac{\partial^2 w}{\partial \phi^2} + \frac{3}{4} \frac{1}{R^2} \left( \frac{E_s}{E_T} - 1 \right) \frac{\partial^6 w}{\partial x^4 \partial \phi^2} \right] = 0 \]
where \( V^4 \) indicates

\[
\left[ \frac{\partial^2}{\partial x^2} + \frac{1}{R^2} \frac{\partial^2}{\partial \phi^2} \right]^2
\]

A solution to this equation can be written:

\[
W = A \sin K \phi \sin \lambda x
\]

(3)

where 

\[
K = N
\]

\[
\lambda = \frac{m \pi}{L}
\]

L is the length of the shell, and

m and n are integers.

This solution satisfies the conditions of simple support at the ends of the cylinder, i.e., that \( w \) and \( \frac{\partial^2 w}{\partial x^2} \) vanish at \( x = 0 \) and \( x = 1 \).

These conditions are not unreasonable for stiffened cylinders since it is likely that the effective rotational restraint will be limited by the formation of plastic regions arising from high bending stresses near stiffeners.

By substituting the solution (3) back into (2), the following characteristic-value equation is obtained:

\[
D \left[ \frac{E_T}{E_s} \left( \frac{K^2}{R^2} + \lambda^2 \right)^4 + \left( 1 - \frac{E_T}{E_s} \right) \lambda^2 \left( \frac{K^2}{R^2} + \lambda^2 \right)^2 \right]
\]

\[
\left[ \frac{3 \lambda^2}{2} + \frac{K^2}{R^2} \right] + \frac{3 \lambda^4}{4} \left( \frac{E_s}{E_T} - 1 \right) \left( \frac{3 \lambda^2}{4} + \frac{K^2}{R^2} \right)
\]
\[ + \frac{E_t}{R^2} \lambda^4 - N_x \left[ \left( \frac{K^2}{R^2} + \lambda^2 \right)^2 \lambda^2 + \frac{3}{4} \left( \frac{E_t}{E_T} - 1 \right) \lambda^6 \right] \\
+ \frac{N \phi}{N_x} \frac{K^2}{R^2} \left\{ \left( \frac{K^2}{R^2} + \lambda^2 \right)^2 + \frac{3}{4} \left( \frac{E_t}{E_T} - 1 \right) \lambda^4 \right\} = 0 \] 

(4)

To simplify:

\[ \phi \equiv \frac{\lambda^2}{\lambda^2 + \frac{K^2}{R^2}} = \frac{1}{1 + \frac{N^2L^2}{m^2\pi^2R^2}} \]

(5)

\[ f \equiv \frac{N_x}{N \phi} \]

\[ c = \left( \frac{E_t}{E_T} - 1 \right) \]

Now since: \[ N_x = \sigma \cdot t = K_x \cdot pt \]

the equation is then rearranged so that an expression for \( P_p \), the plastic buckling pressure, is obtained:

\[
P_p = fD\lambda^2 \frac{E_t}{E_s} \left[ 1 + c \phi \left\{ 1 + \frac{\phi}{2} + \frac{3}{4} \phi^2 \left( 1 - \frac{\phi}{4} \right) \right\} \right] + \frac{E_t f \phi^4}{R^2 \lambda^2} \\
\frac{K_x t \phi \left[ 1 - \phi (1 - f) \right]}{1 + 3 \frac{c \phi^2}{4}}

(6)\]
Equation (6) can now be minimized for $P_p$ with respect to $N/M$ by setting

$$\frac{\partial P_p}{\partial \phi} = 0 \tag{7}$$

$$\phi_p^* = \frac{m^4 \pi^4}{9} \frac{E_T}{E_S} \left( \frac{\sqrt{Rt}}{L} \right)^4 \left[ \frac{1 - 2\phi_p (1-f) + \frac{3}{4}c\phi_p^2}{3 - 2\phi_p (1-f) + \frac{3}{4}c\phi_p^2} \right] \tag{8}$$

and

$$P_p = \frac{4m^2 \pi^2 E_T f}{9\phi_p} \left( \frac{t}{RK_x} \right) \left( \frac{Rt}{L} \right)^2 \left[ \frac{1 + \frac{c\phi_p}{4} (3 + \phi_p + \frac{3}{4}c\phi_p^2)}{3 - 2\phi_p (1-f) + \frac{3}{4}c\phi_p^2} \right] \tag{9}$$

Since $P_p$ in (9) is proportional to $m^2$, $m$ must be equal to one in all cases when $N$ is greater than 0.

After some investigation, an approximate value of $\phi$ was determined

$$\phi \approx 1.23 \frac{\sqrt{Rt}}{L} \tag{10}$$

Equation (9) can be further simplified to

$$P_p = \frac{4\pi^2 E_T f}{9\phi} \left( \frac{t}{RK_x} \right) \left( \frac{\sqrt{Rt}}{L} \right)^2 \left[ \frac{1 + \frac{3c\phi}{4}}{3 - 2\phi (1-f)} \right] \tag{11}$$
A similar analysis for the elastic region yields:

\[
P_e = \frac{\pi^2 E f}{3(1-\nu^2)} \phi \left[ \frac{t}{RK_x} \right] \left( \frac{\sqrt{R t}}{L} \right)^2 \left\{ \frac{1}{3-2\phi(1-f)} \right\}
\]  

(12)

Although Equations (11), (12) define the buckling pressure for the plastic and elastic regions, no solution is given for the inelastic region which lies between these two limiting cases. However, by employing an empirical correction factor wherein Poisson's ratio is regarded as a variable, one can arrive at an expression which reduces to the proper limiting values.

\[
\nu = \frac{1}{2} - \frac{E_s}{E} \left[ \frac{1}{2} - \nu_e \right]
\]  

(13)

Thus rewriting Equation (11),

\[
P_c = \frac{\pi^2 E f}{3(1-\nu^2)} \left[ \frac{t}{RK_x} \right] \left( \frac{\sqrt{R t}}{L} \right)^2 \left[ 1 + \frac{3\phi}{4} \left( \frac{E_s}{E_T} - 1 \right) \right] \left[ \frac{3 - 2\phi(1-f)}{3 - 2\phi(1-f)} \right]
\]  

(14)

2. Inelastic Axisymmetric Buckling Analysis*
The differential equations of equilibrium describing the plastic buckling of a cylindrical shell subjected to external hydrostatic pressure are expressed as Equations (1).

Equations (1) admits to a solution of the form

\[ u = A \cos N\phi \cos \lambda x \]
\[ v = B \sin N\phi \sin \lambda x \]
\[ w = c \cos N\phi \sin \lambda x \]

where \(\lambda = M\pi/L\), \(n\) an integer

This solution describes a buckling mode with \(M\) halfwaves along the cylinder axis and \(2N\) halfwaves around its circumference. It also satisfies boundary conditions of simple support at the ends of the cylinder; i.e., \(w\) and \(\partial^2 w/\partial x^2\) vanish at \(x = 0\) and \(x = L\).

The buckling condition of the shell is represented in Equation (4).

For axisymmetric buckling \(N = 0\) and for \(m = 1\), the minimum value of \(\lambda\) is \(\lambda_{\text{min}} = \pi/L\).

Substituting this value into Equation (4) results in

\[ N_x = \frac{E_t t L^2}{9\pi^2 R^2} \left[ \frac{R^2 t^2 \pi^4 (1 + \frac{3}{4}c)}{L^4} + \frac{9(1+c)}{(1 + \frac{3}{4}c)} \right] \]  

(16)
replacing \( N_x \) with \( \sigma_x t = K_x pt \)

\[
P_p = \frac{E_t L^2}{9K_x \pi^2 R^2} \left[ \frac{\pi^4 R^2 t^2 (1 + \frac{3}{4}c)}{L^4} + \frac{9(1 + c)}{(1 + \frac{3}{4}c)} \right]
\]

(17)

Again using the expression

\[
\nu = \frac{1}{2} - \frac{E_s}{E} \left( \frac{1}{2} - \nu_e \right)
\]

\[
P_c = \frac{E_t L^2}{12K_x (1-\nu^2) \pi^2 R^2} \left[ \frac{\pi^4 R^2 t^2 (1 + \frac{3}{4}c)}{L^4} + \frac{12(1-\nu^2)(1+c)}{(1 + \frac{3}{4}c)} \right]
\]

(18)

3. Web Buckling Under Hydrostatic Loading*

Two modes of buckling will be considered:

(a) Buckling of a Ring

\[ N = 2 \]
(b) Web Buckling under Edge Compression

\[ P_c = (n^2 - 1) \frac{E_T I}{L R_s^3} = \frac{24E_T I_{\text{eff}}}{D_m^3 L} \]  
\( (19) \)

\[ a = 2\pi R_w \]
\[ d = h \]

\[ \sigma_x m^2 + \sigma_y n^2 \left( \frac{a}{d} \right)^2 = \frac{\pi^2 D}{a^2 b} (m^2 + n^2 \frac{a^2}{d^2})^2 \]
\( (20) \)

where

\[ \frac{a}{d} = \frac{2\pi R_w}{h} \gg 1 \]
\[ \sigma_y = 0, \ m > 1, \ n = 1 \]

\[ \sigma_x = \frac{\pi^2 a^2 D}{m^2 b} \left( \frac{m^2}{a^2} + \frac{n^2}{d^2} \right)^2 + \frac{4\pi^2 D}{d^2 b} \] (21)

\[ \sigma_x = 0, \ m = 1, \ n = 1 \]

\[ \sigma_y = \frac{\pi^2 d^2 D}{n^2 b} \left( \frac{m^2}{a^2} + \frac{n^2}{d^2} \right)^2 + \frac{\pi^2 D}{d^2 b} \] (22)

Therefore

\[ \sigma_x + 4\sigma_y = \frac{4\pi^2 D}{d^2 b} \] (23)

and

\[ \sigma_x \equiv \sigma_\phi = K_\phi \rho \]

\[ \sigma_y \equiv \sigma_r = K_r \rho \]

thus

\[ \rho(K_\phi + 4K_r) = \frac{4\pi^2 D}{d^2 b} \]

\[ P_c = \frac{\frac{4\pi^2 D}{d^2 b}}{(K_\phi + 4K_r)} \] (24)
where
\[ d = h \]
\[ D = \frac{Eb}{12(1-v^2)} \]

thus

\[ \frac{P_c}{12(l-v^2)} = \frac{4\pi^2E(b^2/h^2)}{(K_\phi + 4K_r)} \]  (25)

*Reference (2), (4), (8), and (22).*
APPENDIX C

MATERIAL PROPERTIES*

The properties for the four materials utilized in the optimization studies are tabled below. The secant and tangent moduli are defined as

\[ E_s = \frac{\sigma_i}{\epsilon_i} \]

\[ E_T = \frac{d\sigma_i}{d\epsilon_i} \]
(1) **HIGH STRENGTH STEEL** \( (\sigma_y = 160,000 \text{ psi}) \)

<table>
<thead>
<tr>
<th>( \sigma_i / \sigma_y )</th>
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<th>( \sqrt{E_s E_T / E^2} )</th>
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\[ E = 30 \times 10^6 \text{ psi} \]
\[ \nu_E = 0.30 \]
\[ \rho = 490 \text{ pcf} \]

(2) **TITANIUM ALLOY** \( (\sigma_y = 110,000 \text{ psi}) \)

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\[ E = 18 \times 10^6 \text{ psi} \]
\[ \nu_E = 0.30 \]
\[ \rho = 276 \text{ pcf} \]
(3) **ALUMINUM ALLOY** \( (\sigma_y = 65,000) \)

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\( E = 10.8 \times 10^6 \) psi

\( \nu_E = 0.30 \)

\( \rho = 173 \) pcf

(4) **HY80 STEEL** \( (\sigma_y = 80,000 \) psi\)

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\( E = 30 \times 10^6 \) psi

\( \nu_E = 0.30 \)

\( \rho = 490 \) pcf
APPENDIX D

SAMPLE OUTPUT OF COMPUTER PROGRAMS

1. STAS - Stress Analysis Program

2. FMCS - Failure Mode Criteria Program

3. OPTI - Optimization Program
### Stress Analysis of a Web Stiffened Sandwich

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<td>Inside Thickness (in)</td>
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WB = 0.377
KR = 3.071
KP = 83.887

### Stress Analysis of a Web Stiffened Sandwich

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WB = 0.377
KR = 3.071
KP = 83.887

### Stress Analysis of a Web Stiffened Sandwich

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KR = 3.071
KP = 83.887
INSIDE THICKNESS = 1.80
WEB THICKNESS(IN) = 1.80
WEB SPACING(IN) = 16.20

MX00 MPHI00 MX01 MPHI01 MX01 MPHI10 MX11 MPHI11
-37.78 -47.87 -14.75 -40.96 -16.95 -41.91 -35.58 -47.50
FX00 FPHI00 FX01 FPHI01 FX01 FPHI10 FX11 FPHI11

HRW HS HTS BTS IEFF 0
0.050 0.500 5.000 1.000 1369.9 10.800

WB = 0.377

KR = 3.071 KP = 83.887

COMPILE TIME = 7.53 SEC, EXECUTION TIME = 0.86 SEC, OBJECT CODE = 9856 BYT
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**GRAPHICAL DATA FOR FAILURE MODES**

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Stress analysis, buckling analysis and optimum proportions of an isotropic web-stiffened sandwich cylindrical shell under hydrostatic pressure.