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# Extension of Murakami's high-order non-linear solver to multiple roots

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Several one-parameter families of fourth-order methods for finding multiple zeros of non-linear functions are developed. The methods are based on Murakami's fifth-order method (for simple roots) and they require one evaluation of the function and three evaluations of the derivative. The informational efficiency of the methods is the same as the previously developed methods of lower order. For a double root, the method is more efficient than all previously known schemes. All these methods require the knowledge of multiplicity.

**Keywords:** multiple roots; non-linear; high-order; efficiency

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## 1. Introduction

There is a vast literature on the solution of non-linear equations and non-linear systems, see *e.g.* Ostrowski [12], Traub [18], Neta [10] and references therein. Recently several papers by Sharma [15], Sharma and Goyal [16], Homeier [5] and Grau and Diaz Barrero [2] discuss methods for finding simple roots. Here we develop a high-order fixed point type method to approximate a multiple root. There are several methods for computing a zero $\xi$ of multiplicity $m$ of a non-linear equation $f(x) = 0$ (see Neta [10] and Neta and Johnson [11]. Newton's method is only of first order unless it is modified to gain the second order of convergence (see Rall [13] or Schröder [17]). This modification requires a knowledge of the multiplicity. Traub [18] has suggested to use any method for $f^{(m)}(x)$ or $g(x) = f(x)/f'(x)$. Any such method will require higher derivatives than the corresponding one for simple zeros. Moreover, the first one of those methods require the knowledge of the multiplicity $m$. In such a case, there are several other methods developed by Hansen and Patrick [4], Victory and Neta [19], Dong [1] and Neta and Johnson [11]. Since in general one does not know the multiplicity, Traub [18] suggested a way to approximate it during the iteration.

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For example, the quadratically convergent modified Newton's method is

$$x_{n+1} = x_n - m\frac{f_n}{f_n'} \tag{1}$$

and the cubically convergent Halley's method [3] is

$$x_{n+1} = x_n - \frac{f_n}{((m+1)/2m)\,f_n' - (f_n f_n'')/(2f_n')}, \tag{2}$$

where $f_n^{(i)}$ is short for $f^{(i)}(x_n)$. Another third-order method was developed by Victory and Neta [19] and based on King's fifth-order method (for simple roots) [8]

$$\begin{aligned} w_n &= x_n - \frac{f_n}{f_n'}, \\ x_{n+1} &= w_n - \frac{f(w_n)}{f_n'}\frac{f_n + Af(w_n)}{f_n + Bf(w_n)}, \end{aligned} \tag{3}$$

where

$$\begin{aligned} A &= \mu^{2m} - \mu^{m+1} \\ B &= -\frac{\mu^m(m-2)(m-1)+1}{(m-1)^2} \end{aligned} \tag{4}$$

and

$$\mu = \frac{m}{m-1}. \tag{5}$$

Yet two other third-order methods developed by Dong [1], both require the same information and both are based on a family of fourth-order methods (for simple roots) due to Jarratt [6]:

$$x_{n+1} = x_n - u_n - \frac{f_n}{(m/(m-1))^{m+1} f'(x_n - u_n) + (m - m^2 - 1/(m-1))^2 f_n'} \tag{6}$$

$$x_{n+1} = x_n - m/(m+1)u_n - \frac{\frac{m}{m+1} f_n}{(1+(1/m))^m f'(x_n - m/(m+1)u_n) - f_n'} \tag{7}$$

where

$$u_n = \frac{f_n}{f_n'}. \tag{8}$$

Neta and Johnson [11] have developed a fourth-order method based on Jarrat's method [7]. The method in general is given by

$$x_{n+1} = x_n - \frac{f_n}{a_1 f_n' + a_2 f'(y_n) + a_3 f'(\eta_n)} \tag{9}$$

where $u_n$ is given by Equation (8) and

$$\begin{aligned} y_n &= x_n - au_n, \\ v_n &= \frac{f_n}{f'(y_n)}, \\ \eta_n &= x_n - bu_n - cv_n, \end{aligned} \tag{10}$$

where the parameters $a, b, c, a_1, a_2$ and $a_3$ depend on the multiplicity $m$.

Our starting point here is Murakami's two-parameter family of methods [9] given by the iteration

$$x_{n+1} = x_n - a_1 u_n - a_2 w_2(x_n) - a_3 w_3(x_n) - \psi(x_n) \tag{11}$$

where $u_n$ is given by Equation (8) and

$$\begin{aligned} w_2(x_n) &= \frac{f_n}{f'(y_n)}, \quad y_n = x_n - a u_n \\ w_3(x_n) &= \frac{f_n}{f'(z_n)}, \quad z_n = x_n - b u_n - c w_2(x_n) \\ \psi(x_n) &= \frac{f_n}{b_1 f_n' + b_2 f'(y_n)}. \end{aligned} \tag{12}$$

Murakami has shown that this family of methods (for simple roots) is of order 5 [9] if the parameters are chosen appropriately. The method requires one function- and three derivative-evaluation per step. Thus the informational efficiency (see [18]) is 1.25.

## 2. New higher-order scheme

We would like to find the eight parameters $a$, $b$, $c$, $a_1$, $a_2$, $a_3$, $b_1$ and $b_2$ so as to maximize the order of covergence to a root $\xi$ of multiplicity $m$. Let $e_n$, $\hat{e}_n$, $\epsilon_n$ be the errors at the $n$th step, *i.e.*

$$\begin{aligned} e_n &= x_n - \xi, \\ \hat{e}_n &= y_n - \xi, \\ \epsilon_n &= z_n - \xi. \end{aligned} \tag{13}$$

If we expand $f(x_n)$, and $f'(x_n)$ in Taylor series (truncated after the $N$th power, $N > m$) we have

$$f(x_n) = f(x_n - \xi + \xi) = f(\xi + e_n) = \frac{f^{(m)}(\xi)}{m!}\left(e_n^m + \sum_{i=m+1}^{N} A_i e_n^i\right) \tag{14}$$

or

$$f(x_n) = \frac{f^{(m)}(\xi)}{m!} e_n^m \left(1 + \sum_{i=m+1}^{N} B_{i-m} e_n^{i-m}\right) \tag{15}$$

where

$$A_i = \frac{m! f^{(i)}(\xi)}{i! f^{(m)}(\xi)}, \quad i > m \tag{16}$$

$$B_{i-m} = A_i$$

$$f'(x_n) = \frac{f^{(m)}(\xi)}{(m-1)!} e_n^{m-1} \left(1 + \sum_{i=m+1}^{N} \frac{i}{m} B_{i-m} e_n^{i-m}\right). \tag{17}$$

To expand $f'(y_n)$ and $f'(z_n)$ we use some symbolic manipulator, such as Maple [14][1], and find that

$$f'(y_n) = \frac{f^{(m)}(\xi)}{(m-1)!} \hat{e}_n^{m-1} \left(1 + \frac{m+1}{m} B_1 \hat{e}_n + \frac{m+2}{m} B_2 \hat{e}_n^2 + \cdots\right), \tag{18}$$

$$\hat{e}_n = e_n - a u_n = \mu e_n + \frac{a}{m^2} B_1 e_n^2 + \left[\frac{2a}{m^2} B_2 - \frac{a(m+1)}{m^3} B_1^2\right] e_n^3 + \cdots, \tag{19}$$

where

$$\mu = \frac{m-a}{m}. \tag{20}$$

Thus

$$f'(y_n) = \frac{f^{(m)}(\xi)}{(m-1)!} e_n^{m-1} \left(c_0 + c_1 e_n + c_2 e_n^2 + c_3 e_n^3 + \cdots\right), \tag{21}$$

where

$$\begin{aligned}
c_0 &= \mu^{m-1},\\
c_1 &= \mu^{m-2}\frac{(m-a)^2(m+1)+am(m-1)}{m^3}B_1,\\
c_2 &= \frac{\mu^{m-3}}{m^5}\left[am((m-a)^2+m(a+1)^2) - \frac{1}{2}a^2m^2(m+7)\right]B_1^2 + \frac{\mu^{m-3}}{m^5}[(m+2)(m-a)^4\\
&\quad + 2a(m-a)m^2(m-1)]B_2,\\
c_3 &= \mu^m\frac{(m+3)(m-a)^2}{m^3}B_3 - \frac{a\mu^{m-3}}{m^6}[(m^2+3m+2)(a-m)^3 - 2m^2(m+1)(a-m)^2\\
&\quad - 2m^2a(m-1)(m-2)]B_1B_2 + \frac{\mu^{m-4}}{6m^6}a[3(2m^2+4m+1)(a-m)^3 + 2(m+3)(a-m)^2\\
&\quad + a^2m^2(3a+m-18) + am^2(-3m^2+12m+16) - m^2(5m+6)]B_1^3.
\end{aligned} \tag{22}$$

The error in $z_n$ is given by

$$\begin{aligned}
\epsilon_n = e_n - bu_n - cw_2(x_n) &= \left(1 - \frac{b}{m} - \frac{c}{m}\mu^{1-m}\right)e_n\\
&\quad + \left(\frac{b}{m^2} + \frac{c}{m^2}\mu^{-m}\left[\mu^2 - \frac{a(1-a)}{m}\right]\right)B_1e_n^2 + \left[\frac{2b}{m^2} - \left(\frac{\alpha_1}{m^6}c\mu^{-m-1}\right)B_2\right.\\
&\quad \left. - \left(\frac{\alpha_2}{2m^7}c\mu^{-m-1} + \frac{b(m+1)}{m^3}\right)B_1^2\right]e_n^3 + \cdots,
\end{aligned} \tag{23}$$

where

$$\begin{aligned}
\alpha_1 &= -a^4(m+2) + 4a^3m(m+2) - a^2m^2(3m+14) + 2m^3(5a-m)\\
\alpha_2 &= 2a^4m(m+2) - 4a^3m^2(m+4) - 12a^3m + (3m+19)m^3a^2 + 22a^2m^2\\
&\quad - 2am^3(5m+7) + 2a^4 + 2m^5 + 2m^4.
\end{aligned}$$

Now expand $f'(z_n)$ in terms of $e_n$. To this end, we expand in terms of $\epsilon_n$ and substitute for $\epsilon_n$ from Equation (23)

$$\begin{aligned}
f'(z_n) &= \frac{f^{(m)}(\xi)}{(m-1)!}\epsilon_n^{m-1}\left(1 + \frac{m+1}{m}B_1\epsilon_n + \frac{m+2}{m}B_2\epsilon_n^2 + \cdots\right)\\
&= \frac{f^{(m)}(\xi)}{(m-1)!}e_n^{m-1}\left(d_0 + d_1e_n + d_2e_n^2 + d_3e_n^3 + \cdots\right),
\end{aligned} \tag{24}$$

where

$$\begin{aligned} d_0 &= \lambda^{m-1}, \\ d_1 &= \lambda^{m-2}\,\frac{\beta_0+\beta_1 c\mu^{-m}+\beta_2 c^2\mu^{-2m}+\beta_3 bc\mu^{-m}}{m^5}\,B_1, \\ d_2 &= \frac{\lambda^{m-3}}{2m^{10}\mu^{3m+1}}[-(D_1^0+D_1^1\mu^m+D_1^2\mu^{2m}+D_1^3\mu^{3m})B_1^2 \\ &\quad +(D_2^{-1}\mu^{-m}+D_2^0+D_2^1\mu^m+D_2^2\mu^{2m}+D_2^3\mu^{3m})B_2], \end{aligned} \tag{25}$$

and

$$\begin{aligned} \lambda &= \frac{m(m-b)-c(m-a)\mu^{-m}}{m^2}, \\ \beta_0 &= (m^2b(b-m)+m^4\mu^{2m})(m+1), \\ \beta_1 &= m(a^2(m^2-1)+ma(5-m)-m^2(m+3)), \\ \beta_2 &= (m-a)^2(m+1), \\ \beta_3 &= 2m(m+1)(m-a). \end{aligned} \tag{26}$$

The $D_i^j$ are complicated expressions and will be given in the Appendix. Now substitute Equations (15), (17), (21) and (24) into Equations (8), (12) and expand the quotients in Taylor series; then substituting all these into Equation (11), we get

$$e_{n+1} = C_1^1 e_n + C_2^1 B_1 e_n^2 + (C_3^1 B_1^2 + C_3^2 B_2)e_n^3 + (C_4^1 B_1^3 + C_4^2 B_1 B_2 + C_4^3 B_3)e_n^4 + \cdots, \tag{27}$$

where the coefficients $C_i^j$ depend on the parameters $a, b, c, b_1, b_2, a_1, a_2$ and $a_3$. Because of the complexity, we have taken $a = m/2$, $b_2 = 1 - 2^{m-1}b_1$ and either $b = 0$ or $c = 0$. Thus, we reduced the number of parameters to five and as a consequence we were unable to get fifth-order methods. The results for $m = 2$, $m = 3$ and $m = 4$ are given in Table 1.

To summarize, we managed to obtain a family of fourth-order methods requiring one function- and three derivative-evaluation per step. The informational efficiency, $E = p/d$, of these methods

Table 1. The parameters for various values of $m$.

| $m$ | 2 | 2 | 3 | 3 | 4 | 4 |
|---|---|---|---|---|---|---|
| $a$ | 1 | 1 | 3/2 | 3/2 | 2 | 2 |
| $b$ | 0 | Free | 0 | 0.9415780151 | 0 | 11.9151259843 |
| $c$ | Free | 0 | 0.2353945038 | 0 | 1.9640446368 | 0 |
| $b_1$ | 1 | 1 | free | free | 0.05 | 0.0625 |
| $b_2$ | $-1$ | $-1$ | $1-4b_1$ | $1-4b_1$ | 0.0268934369 | 0.5 |
| $a_1$ | $-6$ | $-6$ | $-2.5128989321-16b_1$ | $-10.571320917-16b_1$ | $-7.49156894$ | 5.6116821612 |
| $a_2$ | 3 | 3 | $-1.8238807632+4b_1$ | $0.1907247330+4b_1$ | $-0.91067191$ | $-1.2089575039$ |
| $a_3$ | 0 | 0 | 4.1469082443 | 4.1469082443 | $-0.92646960$ | $-0.4647127230$ |
| $C_4^1$ | 15/32 | 15/32 | $-6.1027059066$ | $-6.1836740792$ | $-1.35078537$ | $-1.0152077055$ |
| $C_4^2$ | $-1/2$ | $-1/2$ | 9.4693139272 | 9.5300400567 | 2.141639816 | 1.5300422793 |
| $C_4^3$ | 1/8 | 1/8 | $-3.4826758270$ | $-3.4826758270$ | $-0.822966734$ | $-0.5494657588$ |

Table 2. Comparison of methods for multiple roots.

| Method | $f$ | $f'$ | $f''$ | $p$ | $d$ | $E = p/d$ | $I = p^{1/d}$ |
|---|---|---|---|---|---|---|---|
| Schröder | 1 | 1 | 0 | 2 | 2 | 1 | 1.4142 |
| Hansen and Patrick | 1 | 1 | 1 | 3 | 3 | 1 | 1.442 |
| Halley | 1 | 1 | 1 | 3 | 3 | 1 | 1.442 |
| Victory and Neta | 2 | 1 | 0 | 3 | 3 | 1 | 1.442 |
| Dong | 1 | 2 | 0 | 3 | 3 | 1 | 1.442 |
| Neta and Johnson | 1 | 3 | 0 | 4 | 4 | 1 | 1.4142 |
| Neta and Johnson, $m = 2$ | 1 | 2 | 0 | 4 | 3 | 1.3333 | 1.5874 |
| Neta | 1 | 3 | 0 | 4 | 4 | 1 | 1.4142 |
| Neta, $m = 2$ | 1 | 2 | 0 | 4 | 3 | 1.3333 | 1.45874 |

is 1, as all the aforementioned methods for multiple roots. The efficiency index, $I = p^{1/d}$, is 1.4142 which is lower than the index for those third-order methods. For $m = 2$, we found that $a_3 = 0$ and thus we need one less derivative. This happened also for the methods developed by Neta and Johnson [11]. In this case, the informational efficiency is 4/3 and the efficiency index is 1.5874. Therefore, our method for double roots is more efficient than the Newton's method as modified by Schröder. If the cost of evaluating the first derivative is lower than that of evaluating the function, then our method, for any multiplicity, will be more efficient than the Newton's method (Equation (1)). These results are given in Table 2. Clearly if the cost of evaluating the derivatives is different than that of evaluating the function, one can make an argument to using the appropriate method for the case at hand.

## 3. Numerical experiments

In all our numerical experiments, we have used the appropriate method with $b = 0$, except for example 3 when we used both schemes. In our first example we took a quadratic polynomial having double roots at $\xi = 1$

$$f(x) = x^2 - 2x + 1. \tag{28}$$

Here we started with $x_0 = 0$ and the root was found in one iteration. The modified Newton method (Equation (1)) converged fast and Newton's method required 10 iterations to get as close as possible to $10^{-7}$. In the second example we took a polynomial having two double roots at $\xi = \pm 1$

$$f(x) = x^4 - 2x^2 + 1. \tag{29}$$

Starting at $x_0 = 0.8$ or $x_0 = 0.6$ our method converged in two iterations. The results are given in Table 3.

Similar results were obtained when starting at $x_0 = -0.8$ to converge to $\xi = -1$. For comparison, we have tried the modified Newton. Using $x_0 = 0.6$ we required four iterations to achieve $10^{-9}$ accuracy.

Table 3. Results for Example 2. $f(x)$ is given by Equation (29).

| $n$ | $x$ | $f$ | $x$ | $f$ |
|---|---|---|---|---|
| 0 | 0.8 | 0.1296 | 0.6 | 0.4096 |
| 1 | 1.00100728 | 0.4062524998 (−5) | 1.03262653 | 0.004398017 |
| 2 | 1.00000000 | 0 | 1.00000036 | 0.5060180 (−12) |

Table 4. Results for Example 3. The first three columns use the scheme with $b = 0$ and the last 3 use $c = 0$. $f(x)$ is given by Equation (30).

| $n$ | $x$ | $f$ | $n$ | $x$ | $f$ |
|---|---|---|---|---|---|
| 0 | 0 | −6 | 0 | 0 | −6 |
| 1 | 0.989582711 | −0.22964188 (−5) | 1 | 0.985370624 | −0.64000004 (−5) |
| 2 | 0.999999994 | 1 (−18) | 2 | 0.999999974 | 0 |

Table 5. Results for Example 4. $f(x)$ is given by Equation (31).

| $n$ | $x$ | $f$ | $x$ | $f$ |
|---|---|---|---|---|
| 0 | 0.1 | 0.11051709 (−1) | 0.2 | 0.4885611033 (−1) |
| 1 | 0.2069496569 (−4) | 0.428290468 (−9) | 0.286951344 (−3) | 0.8236470507 (−7) |
| 2 | 0.43944 (−19) | 0.193107514 (−38) | 0.162369865 (−14) | 0.2636397306 (−29) |

The next example is a polynomial with triple root at $\xi = 1$

$$f(x) = x^5 - 8x^4 + 24x^3 - 34x^2 + 23x - 6. \tag{30}$$

The iteration starts with $x_0 = 0$ and the results are summarized in Table 4. The first three columns use the scheme with $b = 0$ and the last three columns use $c = 0$.

Another example with double root at $\xi = 0$ is

$$f(x) = x^2 e^x. \tag{31}$$

Starting at $x_0 = 0.1$ or even $x = 0.2$ our method converged in two iterations. The results are given in Table 5.

The next example having a double root at $\xi = 1$ is

$$f(x) = 3x^4 + 8x^3 - 6x^2 - 24x + 19. \tag{32}$$

Now we started with $x_0 = 0.5$ and the results are summarized in Table 6.

The last example having a root at $\xi = 1$ with multiplicity $m = 4$ and a simple root at $\xi = -1$, *i.e.*

$$f(x) = x^5 - 3x^4 + 2x^3 + 2x^2 - 3x + 1. \tag{33}$$

Now we started with $x_0 = 0.01$ and the results are summarized in Table 7. This is the only case where we needed more than two iterations to converge.

Table 6. Results for Example 5. $f(x)$ is given by Equation (32).

| $n$ | $x$ | $f$ |
|---|---|---|
| 0 | 0.5 | 6.6875 |
| 1 | 1.00806166565 | 0.235014761 (−2) |
| 2 | 1.00000000024 | 0.2 (−17) |

Table 7. Results for Example 6. $f(x)$ is given by Equation (3).

| $n$ | $x$ | $f$ |
|---|---|---|
| 0 | 0.01 | 0.9702019701 |
| 1 | 0.090514708167 | 0.905147081668 (−1) |
| 2 | 0.562284899208 | 0.573490665693 (−1) |
| 3 | 0.993019776872 | 0.47313908958 (−8) |
| 4 | 0.999999999699 | 0 |

## 4. Conclusions

We have extended Murakami's method to obtain non-simple zeros. We have developed a one-parameter family of fourth-order methods for various values of the multiplicity. The methods listed are not the only solution to the system of equations and we only listed a representative scheme. The numerical experiments demonstrate the rapid convergence of our method. Because of the complexity of the symbolic manipulation, we had to assign certain values to some of the parameters and were unable to achieve fifth order.



## Note

1. 'Maple is a system for mathematical computation symbolic, numerical and graphical' [14]. For example, the command 'rfp:=series(1/fp1,e,6);' expands the reciprocal of the function $fp1$ into a power series in the variable $e$ keeping terms up to $O(e^6)$. Maple can convert that expansion into a polynomial by using the command 'rfp:=convert(rfp,polynom):'

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## Appendix

Here we list the coefficients $D_i^j$ in the expression for $d_2$ in Equation (25).

$$D_1^0 = -26c^3a^2m^4 + 32m^4c^3a^3 + 26c^3a^3m^3 - 18c^3a^4m^3 - 28m^5c^3a^2 - 2m^6c^3 - 2m^6c^3a^2 + 6m^5c^3a^3 - 2m^7c^3 + 2c^3a^5m^3 + 12m^6c^3a - 12c^3a^4m^2 + 4c^3a^5m^2 - 6c^3a^4m^4 + 12c^3am^5 + 2c^3a^5m.$$

$$
\begin{aligned}
D_1^1 = {} & -37m^6c^2a - 2m^5c^2a - m^7c^2a - 4m^6c^2a^2b + 7m^6c^2a^3 - m^6c^2a^2 - m^6c^2a^4 \\
& + 72m^5c^2a^2 - 8c^2a^4m^3b + 26m^6bac^2 - 38bc^2a^2m^4 - 4c^2a^4m^2b - m^7c^2a^2 \\
& + 8m^5c^2a^3b + m^4c^2a^5 - 6bc^2m^6 + 22bc^2a^3m^3 - 7m^5c^2a^4 + 2m^4c^2a^2 \\
& + 26bc^2am^5 + m^8c^2 - 42bc^2a^2m^5 - c^2a^5m^3 - 7m^5c^2a^3 + 7m^7c^2 - c^2a^5m^2 \\
& + 9m^4c^2a^4 - 4m^4c^2a^4b + 15c^2a^4m^3 + 30bc^2a^3m^4 + m^5c^2a^5 - 56c^2a^3m^4 - 6m^7c^2b. \\
D_1^2 = {} & -10ca^3m^4b + 2ca^4m^2b - 2m^6ca^2b^2 - 5m^6ca^2 - 8m^3bca^3 - 2m^7c + 2ca^4m^3b \\
& + 6m^7ca^2 + 14m^7bc - 2m^7bca + 2m^5ca^3b^2 + 8m^4bca^2 + 16b^2cam^5 + 14m^6ca \\
& - 2m^8ca + 4m^7ca - m^7ca^2b - 2ca^4m^3 - 2ca^4m^4b + 2m^6ca^4 - 6b^2m^6c - 2m^8c \\
& - 8m^6ca^3 + 2b^2ca^3m^3 - 22ca^2m^5 + 6ca^3m^5 - 2m^7ca^3 + m^8ca^2 + 2m^6ca^3b \\
& - 44m^6bca + 16b^2am^6c + 2m^8bc - 2ca^4m^4 + 4ca^3m^4b^2 + 12ca^3m^4 - 12b^2a^2m^4c \\
& + 37m^5ba^2c - 2m^5ca^4b - 14b^2a^2m^5c - 2m^5bca - 6m^7cb^2 + 4m^6ba^2c + 2ca^4m^5. \\
D_1^3 = {} & 2b^3am^5 - 7m^6b^2a + 2b^3am^6 + 7m^7b^2 + 2m^7ba - 2m^8b + m^8b^2 - 2b^3m^6 \\
& - 2m^7b^3 - 2bm^7 - m^7b^2a + 2bam^6. \\
D_2^{-1} = {} & 20c^4m^4a + 40c^4a^3m^2 - 40c^4a^2m^3 - 20m^4c^4a^2 + 2c^4a^5m + 4c^4a^5 + 10m^5c^4a \\
& - 2m^6c^4 - 10c^4a^4m^2 - 20c^4a^4m + 20c^4a^3m^3 - 4c^4m^5. \\
D_2^0 = {} & -64c^3a^3m^3 - 64c^3am^5 - 16bc^3m^5 + 96c^3a^2m^4 - 8m^6bc^3 + 8c^3a^4m^3 - 32m^6c^3a \\
& - 32m^4c^3a^3 + 16c^3a^4m^2 + 64bc^3a^3m^2 + 48m^5c^3a^2 - 16bc^3a^4m - 96bc^3a^2m^3 \\
& + 64bc^3am^4 + 32m^5bc^3a - 8bc^3a^4m^2 + 32bc^3a^3m^3 + 16m^6c^3 - 48bc^3a^2m^4. \\
D_2^1 = {} & -48bc^2a^3m^3 - 14m^6c^2a^3 - 144bc^2am^5 - 72m^6bac^2 + 12b^2a^3m^3c^2 - 36b^2a^2m^4c^2 \\
& + 24b^2a^3m^2c^2 - 72b^2c^2a^2m^3 - 24bc^2a^3m^4 - 20c^2a^4m^3 + 4c^2a^5m^2 + 10m^4c^2a^4 \\
& + 68c^2a^3m^4 + 96m^6c^2a - 2m^4c^2a^5 + 12m^7c^2a + 24m^7c^2b + 36b^2am^5c^2 \\
& + 72bc^2a^2m^5 - 12m^6b^2c^2 + 6m^7c^2a^2 + 144bc^2a^2m^4 \\
& - 120m^5c^2a^2 + 72b^2c^2am^4 - 8m^8c^2 - 18m^5c^2a^3 - 24b^2m^5c^2 - 28m^7c^2 + 10m^5c^2a^4 \\
& + 6m^6c^2a^2 + 48bc^2m^6 - 2c^2a^5m^3. \\
D_2^2 = {} & -14m^7ca^2 + 4m^8ca + 20m^8c + 6m^7ca^2b - 80m^5ba^2c - 96b^2cam^5 + 16b^3am^5c \\
& - 8m^6ca^3b - 8m^5ca^3b + 8m^7c^3 + 2m^5ca^4b - 16ca^3m^5 + 16ca^3m^4b - 8b^3a^2m^4c \\
& - 16b^3m^5c - 16m^8bc + 24b^2a^2m^5c + 124m^6bca - 4ca^4m^3b - 8m^6b^3c + 32b^3cam^4 \\
& + 48b^2a^2m^4c + 4ca^4m^4 + 24m^7cb^2 - 16b^3a^2m^3c + 8m^6ca^3 - 2ca^4m^5 + 2m^6ba^2c \\
& + 44m^6ca^2 + 2ca^4m^4b - 48b^2am^6c + 20m^7bca - 6m^8ca^2 + 4m^9c - 56m^7bc \\
& - 2m^6ca^4 + 8m^7ca^3 - 52m^7ca + 48b^2m^6c. \\
D_2^3 = {} & -20m^7ba + 2b^4am^5 - 4m^8ba + 4m^9b - 4m^9 - 2m^{10} - 4b^4m^5 + 20m^8b - 16b^3am^5 \\
& + 8m^7b^2a - 8b^3am^6 + 4b^4am^4 + 28m^6b^2a - 2m^6b^4 - 8m^8b^2 + 8m^7b^3 - 28m^7b^2 \\
& + 2m^9a + 16b^3m^6 + 4m^8a.
\end{aligned}
$$